



# Colloidal gels

#### Multiscale approaches to their structure and rheology

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#### Overview and schedule

Course goal: Highlight the key physical phenomena driving colloidal gelation and gel properties and discuss experimental methods to study their structure and rheology, with an emphasis on methods that enable an understanding of the connections between different length scales.

#### LECTURE 1

- 1. Colloidal suspensions
  - a. Characteristics of the colloidal domain
- 2. Colloidal interactions
  - a. van der Waals interactions
  - b. Electrostatic interactions
    - i. polar liquids
    - ii. non-polar liquids
  - c. Grafted and adsorbed polymers and surfactants
  - d. Model interactions

#### Reading:

Israelachvili, chapter X

#### **LECTURE 2**

- 3. Colloid motion
  - a. Mobility and resistance
  - b. Hydrodynamic interactions
  - c. Brownian motion
- 4. Equilibrium
  - a. Equilibrium phases of repulsive colloids
  - b. Equilibrium phases of attractive colloids
  - c. Non-equilibrium states glasses and gels

#### Reading:

Furst and Squires, chapters 2, 3

#### LECTURE 3

- 5. Strongly flocculated suspensions
  - a. Colloidal interactions of strongly flocculated suspensions
  - b. Kinetics of flocculation
  - c. Quiescent structure of flocculated suspensions fractal gels
- 6. Gel rheology
  - a. Elastic modulus
  - b. Yielding

#### LECTURE 4

- 7. Weakly flocculated suspensions
  - a. Colloidal interactions of weakly flocculated suspensions
  - b. Structure of weakly flocculated suspensions and gels
- 8. Sedimentation, settling, and redispersion
  - a. Delayed collapse
  - b. Poroelastic compaction

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#### Overview

#### **LECTURE 5**

- 9. Methods for characterizing micro-scale structure and dynamics
  - a. Scattering
  - b. Microscopy
  - c. Dynamics of fractal gels
  - d. Dynamics of dense suspensions
- 10. Micromanipulation and particle interaction measurements
  - a. Surface and colloidal force measurements
  - b. Laser tweezers
  - c. Mechanics of colloidal bond rupture
  - d. Bond mechanics of strongly flocculated suspensions

#### **LECTURE 6**

- 11. Multiscale models of gel rheology
  - a. Strong gels and the fractal model
  - b. Failure of fractal models and motivation for new multiscale models
  - c. Cluster gels
- 12. Gel structure under shear
  - a. Compaction
  - b. Startup of steady shear
  - c. Thixotropy

# Colloid science

# What is a colloid?



polystrene latex, 1.8µm diameter



The IUPAC definition of a colloid is a ``state of subdivision such that... particles dispersed in a medium have at least one dimension between approximately 1 nm and 1  $\mu$ m..."

A colloid is *large enough* to be much larger than the atoms or molecules around it (or inside it), but *small enough* that it exhibits some properties of the molecular scale (e.g. Brownian motion).

12 nearest neighbors in an fcc crystal (0.74)

$$\frac{4}{3}\pi(2a)^3 - \frac{4}{3}\pi a^3 = \frac{28}{3}\pi a^3$$

How many H<sub>2</sub>O molecules if  $2a = 1 \ \mu \text{m}$ ?

$$\bar{V} = 3 \times 10^{-29} \text{ m}^3$$
  
 $V/\bar{V} = 1.2 \times 10^{11} \text{ molecules}$   
 $\bar{V}^{-1/3} \sim 3.1 \text{\AA}$ 

# What is a colloid?



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Continuum mechanics (inside and outside)
Continuum electrodynamic properties (dielectric constant, index of refraction, etc.)



## Brownian motion



## Time scales

		2a = 100  nm	$1\mu{ m m}$	$10\mu{ m m}$
Vorticity	$\tau_{\eta} = \frac{a^2}{\nu} = \frac{a^2 \eta_s}{\rho_s}$	2.5 ns	250 ns	
Inertial	$\tau_I = \frac{2\rho_p}{9\rho_s}\tau_\eta$	0.6 ns	60 ns	
Brownian	$\tau_B = \frac{\pi \eta a^3}{kT}$	0.1 ms	100 ms	100 s
$ u \approx 10^6 \mathrm{m^2/s} $ $ kT = 4.1 \times 10^{-21} \mathrm{J} $				

# Force and energy scales

 $kT = 4.1 \times 10^{-21} \text{ J}$ 

(4 zeptajoules)

		2a = 1  nm	$1\mu{ m m}$
Brownian force	$F_B = \frac{kT}{a}$	10 <sup>-12</sup> N	10 <sup>-15</sup> N
Brownian stress	$\sigma_B = \frac{kT}{a^3}$	10 <sup>3</sup> Pa	10 <sup>-6</sup> Pa

# A colloid?





# Examples of colloids





Rough / smooth Monodisperse / polydisperse Isotropic / anisotropic Uniform / patchy Prolate / oblate

# What is a colloidal gel?

- An elastic solid with a yield stress composed of...
- Percolated microstructure
- and (incompressible) solvent

## Colloidal gels

K.A. Whitaker et al., Nature Commun. 10, 2237 (2019)





# Fundamental and economic significance

- Colloidal suspensions with attractive interparticle interactions include colloid-polymer mixtures, colloids with screened electrostatic interactions, protein solutions at low temperatures.
- Exhibit a wide range of equilibrium and non-equilibrium behavior
- This affects the stability and shelf-life of consumer products, foods, agrochemicals, pharmaceuticals, coatings, paints, etc.
- Generalizable / predictable behavior to aid development and engineering?

#### Examples of equilibrium and non-equilibrium behavior

Phase separation



Stable, homogeneous dispersion



Thickening / gelation



Consolidation / sedimentation



# Colloidal interactions and stability

## Colloidal interactions

Israelachvili, *Intermolecular and Surface Forces*, 3rd ed., Academic Press, New York, 2007.



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#### van der Waals interactions



Related to the polarizability of molecules

$$\begin{array}{ll} \text{Iarizing} & u_{\mathrm{ind}} = ql = \alpha E \\ & E_r = -2\alpha E/4\pi\epsilon\epsilon_0 r^3 = -2\alpha(ze)/(4\pi\epsilon\epsilon_0)^2 r^5 \\ & F(r) = (ze)E_r = -2\alpha(ze)^2/(4\pi\epsilon\epsilon_0)^2 r^5 \\ & w(r) = -\int_{\infty}^0 Fdr = -\alpha(ze)^2/2(4\pi\epsilon\epsilon_0)^2 r^4 \end{array}$$

Keesom interaction angle-averaged dipole-dipole  $u_1 \checkmark \checkmark u_2$   $w(r) = \frac{-u_1^2 u_2^2}{3(4\pi\epsilon\epsilon_0)^2 kT} \left(\frac{1}{r^6}\right)$ Debye interaction rotating dipole-induced dipole  $u_1 \checkmark \checkmark \alpha_0$   $w(r) = \frac{-u_1^2 \alpha_0}{(4\pi\epsilon\epsilon_0)^2} \left(\frac{1}{r^6}\right)$ London interaction induced dipole induced dipole  $\alpha_0 \diamond r \diamond \alpha_0$   $w(r) = \frac{-3\alpha_0^2 h\nu}{4(4\pi\epsilon_0)^2} \left(\frac{1}{r^6}\right)$ 

## London dispersion interaction

$$\alpha_0 \odot \stackrel{r}{\longrightarrow} \alpha_0 \qquad w(r) = \frac{-3\alpha_0^2 h\nu}{4(4\pi\epsilon_0)^2} \left(\frac{1}{r^6}\right)$$



momentary dipole  $\mathcal{O}(10^{-16})\,\mathrm{s}$ 



 $\mathbf{E} = \frac{1}{4\pi\epsilon_0 r^3} \left( 3\hat{\mathbf{x}}\hat{\mathbf{x}} - \mathbf{I} \right) \cdot \mathbf{u}$ 

$$\nu \sim 3.3 \times 10^{15} \,\mathrm{s}^{-1}$$
$$I = h\nu$$

First ionization potential of Bohr atom based on its "orbital frequency"

- "Long-range" interaction 0.2-10nm
- Usually attractive, special cases repulsive
- Brings together molecules and weakly aligns
- Not pairwise additive

How strong? Consider two atoms with

$$\frac{\alpha_0}{4\pi\epsilon_0} = 1.5 \times 10^{-30} \text{ m}^3$$
$$I = h\nu = 2 \times 10^{-18} \text{ J}$$
$$r \sim \sigma \approx 0.3 \text{ nm}$$

$$w(\sigma) = -4.6 \times 10^{-21} \,\mathrm{J}$$

 $\approx 1 \,\mathrm{kT}$ 

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#### van der Waals interactions



**FIGURE 11.2** Methods of summing (integrating) the interaction energies between molecules in condensed phases to obtain the interaction energies between macroscopic bodies. (a) Molecule near a flat surface or "wall." (b) Spherical particle near a wall ( $R \gg D$ ). (c) Two planar surfaces ( $l \gg D$ ).

#### van der Waals attraction

Geometry of bodies with _ surfaces <i>D</i> apart ( <i>D</i> « <i>R</i> )		Van der Waals Interaction*		
		Energy, W	Force, $F = -dW/dD$	
Two atoms or small molecules	TWO ATOMS & SMALL MOLECULES	-C/r <sup>6</sup>	-6C/r <sup>7</sup>	
Two flat surfaces (per unit area)	TWO FLAT SUPPACES	$W_{\rm flat} = -A/12\pi D^2$	$-A/6\pi D^3$	
Two spheres or macromolecules of radii R <sub>1</sub> and R <sub>2</sub>	TWO SPHERES $\begin{array}{c} & & \\ & & $	$\frac{-A}{6D} \left( \frac{R_1 R_2}{R_1 + R_2} \right)$	$\frac{-A}{6D^2} \left( \frac{R_1 R_2}{R_1 + R_2} \right)$ Also $F = 2\pi \left( \frac{R_1 R_2}{R_1 + R_2} \right) W_{\text{flat}}$	
Sphere or macro- molecule of radius <i>R</i> near a flat surface	SPHERE ON FLAT	−AR/6D	$-AR/6D^2$ Also $F = 2\pi RW_{flat}$	
Two parallel cylinders or rods of radii $R_1$ and $R_2$ (per unit length)	TWO MAALLEL CYLINDERS	$\frac{-A}{12\sqrt{2}D^{3/2}} \left(\frac{R_1R_2}{R_1+R_2}\right)^{1/2}$	$\frac{-A}{8\sqrt{2}D^{5/2}} \left(\frac{R_1R_2}{R_1+R_2}\right)^{1/2}$	
Cylinder of radius R near a flat surface (per unit length)	CYLINDER ON FLAT	$\frac{-A\sqrt{R}}{12\sqrt{2}D^{3/2}}$	$\frac{-A\sqrt{R}}{8\sqrt{2}D^{5/2}}$	
Two cylinders or filaments of radii R <sub>1</sub> and R <sub>2</sub> crossed at 90°	CROSSED CYLINDERS $R_1$ $R_2$ $R_3$ $R_4$ $R_5$ $R_2 \gg 0$	$\frac{-A\sqrt{R_1R_2}}{6D}$	$\frac{-A\sqrt{R_1R_2}}{6D^2}$ Also $F = 2\pi\sqrt{R_1R_2}W_{\text{flat}}$	

**FIGURE 13.1** Van der Waals interaction energy *W* and force *F* between macroscopic bodies of different geometries in terms of their Hamaker Constant, *A*. Negative *F* implies attraction (A positive); positive *F* means repulsion (A negative). The Hamaker constant *A* is defined as  $A = \pi^2 C \rho_1 \rho_2$  where  $\rho_1$  and  $\rho_2$  are the number of atoms per unit volume in the two bodies and *C* is the coefficient in the atom-atom pair potential (top row). More rigorous methods for calculating Hamaker constants are described in Sections 13.4 and 13.5. For other geometrics, see Parsegian (2006).  $A_H$  Hamaker constant

$$w(r) = -C/r^6$$

$$A_H = \pi^2 C \rho_1 \rho_2$$

$$A_H = (0.4-4) \times 10^{-19} \,\mathrm{J}$$

#### Interaction between spheres (more rigorous)

$$\Phi_{\rm vdw}(h) = -\frac{A_H}{12} \left\{ \left( \frac{4a^2}{4ah + h^2} \right) + \left( \frac{2a}{2a + h} \right)^2 + 2\ln\left[ 1 - \left( \frac{2a}{2a + h} \right) \right] \right\}$$

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#### Hamaker constants

			Absorption	Hamaker Constant A ( $10^{-20}$ J)		
Medium	Dielectric Constant $\varepsilon$	Refractive Index <i>n</i>	Frequency v <sub>e</sub> (10 <sup>15</sup> s <sup>-1</sup> )	Eq. (13.16) ε <sub>3</sub> = 1	Exact solutions <sup>a</sup>	Experiment <sup>b</sup>
Liquid He	1.057	1.028	5.9	0.057		
Water	80	1.333	3.0	3.7	3.7–5.5	
<i>n</i> -Pentane ( $C_5H_{12}$ )	1.84	1.349	3.0	3.8	3.75	
<i>n</i> -Octane	1.95	1.387	3.0	4.5	4.5	
<i>n</i> -Dodecane	2.01	1.411	3.0	5.0	5.0	
<i>n</i> -Hexadecane	2.05	1.423	2.9	5.1	5.2	
Hydrocarbon (crystal)	2.25	1.50	3.0	7.1		10
Diamond	5.66	2.375	2.6	28.9	29.6	
Cyclohexane (C <sub>6</sub> H <sub>12</sub> )	2.03	1.426	2.9	5.2		
Benzene ( $C_6H_6$ )	2.28	1.501	2.1	5.0		
Carbon tetrachloride ( $CCl_4$ )	2.24	1.460	2.7	5.5		
Acetone (CH <sub>3</sub> ) <sub>2</sub> CO	21	1.359	2.9	4.1		
Ethanol ( $C_2H_5OH$ )	26	1.361	3.0	4.2		
Polystyrene	2.55	1.557	2.3	6.5	6.6-7.9	
Polyvinyl chloride	3.2	1.527	2.9	7.5	7.8	
PTFE (Teflon)	2.1	1.359	2.9	3.8	3.8	
Silica (SiO <sub>2</sub> )	3.8	1.448	3.2	6.3	6.5	5–6
Mica	5.4-7.0	1.60	3.0	10	7–10	13.5
CaF <sub>2</sub>	7.4	1.427	3.8	7.0	7.0	
Silicon (Si)	11.6	3.44	0.80	18	19–21	
Silicon nitride (Si <sub>3</sub> N <sub>4</sub> )	8	1.98	2.45	17	17	
Silicon carbide (SiC)	10.2	2.65	1.8	25	25	
$\alpha$ -Alumina sapphire (Al <sub>2</sub> O <sub>2</sub> )	10 1–11 6	1 75	3.2	15	15	
Zirconia $(n-7r\Omega_2)$	18	2 15	2.1	18	20	
Zinc sulfide (ZnS)	8.5	2.26	1.6	16	15-17	
Metals (Au Ag Cu)	∞		3–5	25-40	20-50	
				Eq. (13.19)		

**Table 13.2**Nonretarded Hamaker Constants for Two Identical Media Interacting in<br/>a Vacuum (Inert Air) at Room Temperature

<sup>a</sup>Exact solutions computed by Hough and White (1980), Parsegian and Weiss (1981), H. Christenson (1983, thesis), Velamakanni (1990, thesis), Senden et al., (1995), French et al., (1995), Bergström et al., (1996), Bergström (1997), Parsegian (2006). <sup>b</sup>Experimental values from Israelachvili and Tabor (1972), Derjaguin et al., (1978), and other literature sources.

## **Electrostatic repulsion**



Poisson equation

$$-\epsilon\epsilon_0\nabla^2\psi = ze\rho$$

Boltzmann distribution

$$\rho = \rho_0 e^{-ze\psi/kT}$$

#### Poisson-Boltzmann equation

$$\nabla^2 \psi = -(ze\rho_0/-\epsilon\epsilon_0)e^{-ze\psi/kT}$$

Israelachvili, <u>Intermolecular and Surface Forces</u>, 3rd ed., Academic Press, New York, 2007.

Russel, Saville, and Schowalter. <u>Colloidal Dispersions</u>. Cambridge University Press, New York, 1989.

## Interaction potential between spheres



Constant potential surfaces

$$U_{\rm el}(h) = 2\pi\epsilon\epsilon_0 (kT/ze)^2 a\Psi_s^2 \ln\left[1 + e^{-\kappa h}\right]$$

Debye length

Ionic strength

$$\kappa^{-1} = \left(\frac{kT\epsilon\epsilon_0}{2Ie^2}\right)^{\frac{1}{2}}$$

$$I = \frac{1}{2} \sum_{i=1}^{N} z_i^2 n_{s,i}$$

$$n_{s,i} = 1000 N_A c_{s,i}$$

NaCl (mM)	$\kappa^{-1}$ (nm)
0.1	31
1	9.7
10	3.1
100	0.97
	in water, R.T.

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#### Kinetic stability



Brownian  $au_B = \frac{\pi \eta a^3}{kT}$   $\frac{2a = 100 \text{ nm}}{0.1 \text{ ms}}$   $\frac{1 \,\mu\text{m}}{10 \,\mu\text{m}}$ 

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# Cohesion

When particles come into contact, they are held by strong van der Waals forces





(b)

$$W_{SL} = \frac{1}{2}W_{SS} + \frac{1}{2}W_{LL} - \gamma_{SL}$$

Then

$$W_{SL} = \gamma_S + \gamma_L - \gamma_{SL}$$

Young equation

$$\gamma_S = \gamma_{SL} + \gamma_L \cos \theta$$

Young-Dupré equation

$$W_{SL} = \gamma_L (1 + \cos \theta)$$

Alternately,

$$\gamma_S \approx \frac{A_H}{12\pi D_0^2}$$



## Adhesion between elastic particles JKR Theory

Johnson, Kendall and Roberts, Proc. R. Soc. London, Ser. A 324, 301 (1971).

 $P_0 = 0$   $a_c = \left[\frac{3\pi a^2 W_{SL}}{2K}\right]^{1/3}$ 



Surface adhesion energy:

$$W_{SL} = \gamma_S + \gamma_L - \gamma_{SL}$$

$$W_{SL} = \gamma_L (1 + \cos \theta_0) \approx 93 \mathrm{mJ/m^2}$$
  
Young-Dupré eqn

where  $\theta_0 = 73.7 \pm 0.3^{\circ}$ 

 $\gamma_L \approx 72 - 73 \text{mN/m}$ 

Elastic modulus: 
$$K = \frac{2E}{3(1-\nu^2)}$$
  
 $K \approx 2.5$ GPa <sup>2</sup>

1 Kwok et al. J. Colloid Interface Sci. 206, 44 (1998). 2 Jarvis and Scheiman, J. Phys. Chem. 72, 74 (1968).

#### DMT and JKR theories

Derjaguin, B., Muller, V. & Toporov, Y., J. Colloid Interface Sci. 53, 314–326 (1975) Tabor, D. J. Colloid Interface Sci. 1977, 58, 2--13.



 $2a = 1\mu m, W \sim 100 m J/m^2, h \sim 0.3 n m$ 

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## Energy holding particles together

Johnson, Kendall and Roberts, Proc. R. Soc. London, Ser. A 324, 301 (1971).



We should calculate the force!

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## Steric forces

D. H. Napper. Polymeric Stabilization of Colloidal Dispersions. Academic Press, New York, 1983.

Adsorbed and grafted polymers Homopolymers and co-polymers



For chemically-grafted polymers:

$$\Phi_p = \begin{cases} \Phi_0 \left[ -\ln y - \frac{9}{5}(1-y) + \frac{1}{3}(1-y^3) - \frac{1}{30}(1-y^6) \right], & 0 < y < 1\\ 0 & y \ge 1 \end{cases}$$

$$y = h/2L = (r - 2a)/2L$$

$$\Phi_0 = \left(\frac{\pi L \sigma_p k_B T}{12 N_p l^2}\right) a L^2$$

 $\begin{array}{l} L \mbox{ contour length} \\ N_p \mbox{ degree of polymerization} \\ l \mbox{ segment length} \\ \phi_p \mbox{ surface graft fraction} \end{array}$ 

## **Bridging interactions**

Adsorbed polymer, protein at lower surface coverage





G.J. Fleer, M.A. Cohen Stuart, J.M.H.M. Scheutjens, T. Cosgrove, B. Vincent. Polymers at Interfaces. New York, 1993.

#### **Depletion attraction**

Asakura, S. & Oosawa, F. J Chem Phys 22, 1255–1256 (1954). Vrij, A. Pure Appl. Chem. 48, 471–483 (1976).

