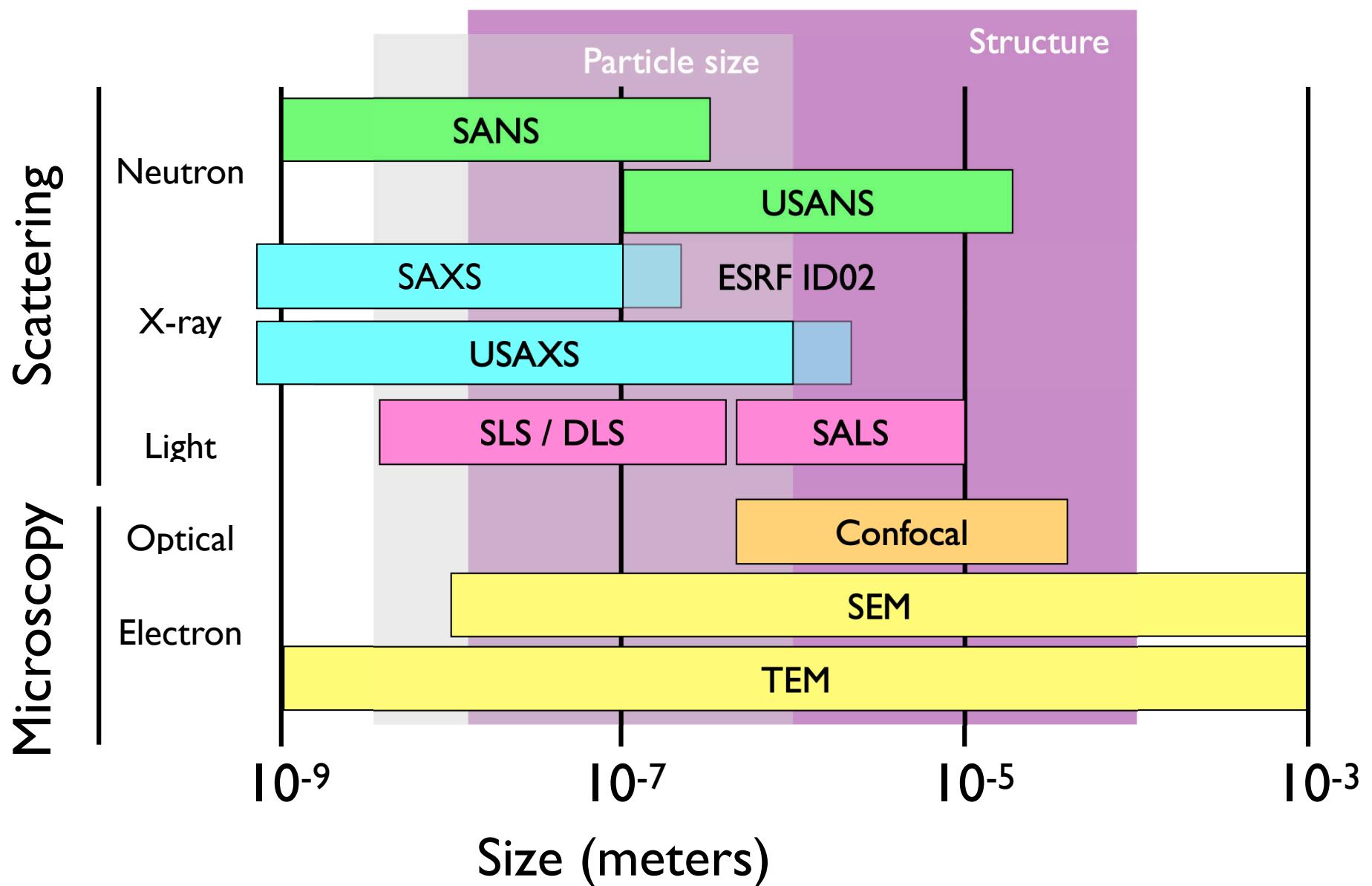


Methods of structural characterization

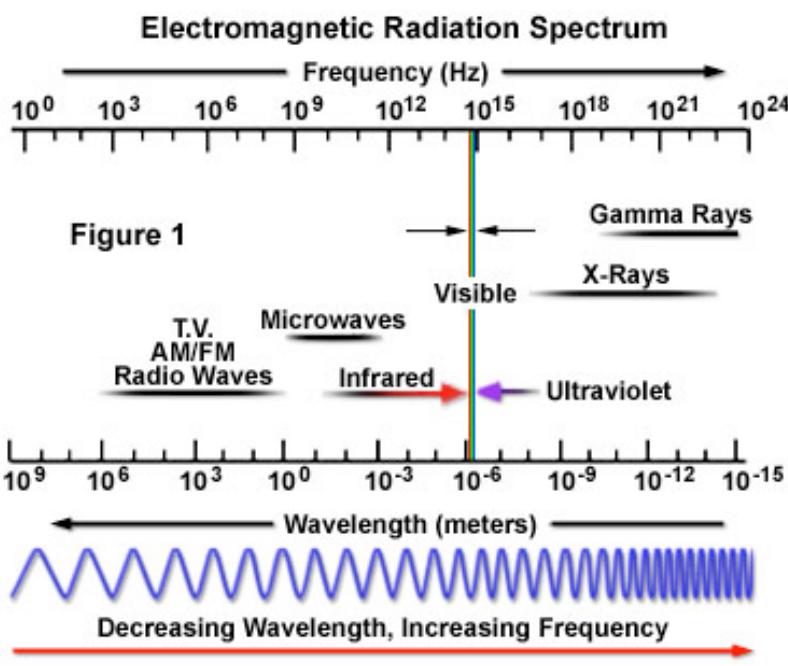
Structure characterization



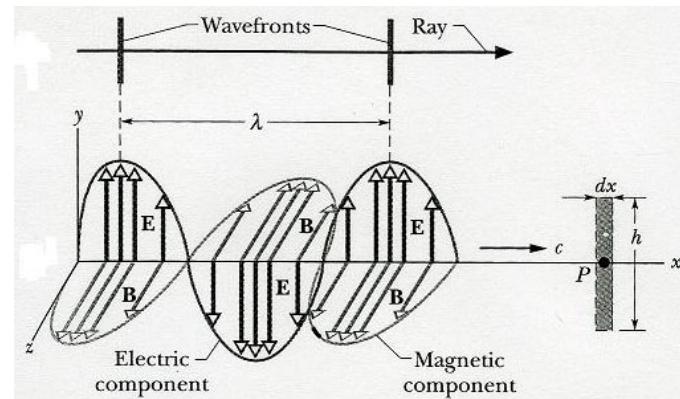
Scattering

Scattering basics

Visible light and x-rays are transverse electromagnetic waves with perpendicular electric (E) and magnetic (B) components.



<http://www.olympusmicro.com/primer/lightandcolor/electromagintro.html>



Halliday, Resnick, Walker Fundamentals of Physics, 1997

$$\mathbf{E}(z, t) = E_x \cos(kz - \omega t) \mathbf{e}_x + E_y \cos(kz - \omega t) \mathbf{e}_y$$

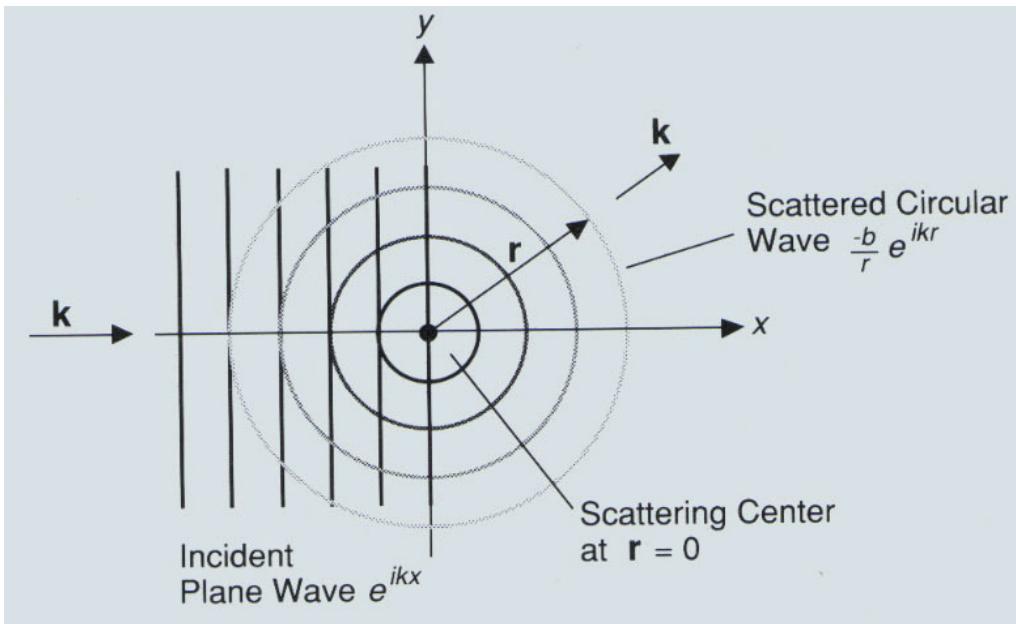
Cold neutrons (5×10^{-5} - 0.025 eV):

$$\lambda_N = 0.2 - 2 \text{ nm}$$

$$\lambda = h/mv$$

Scattering basics

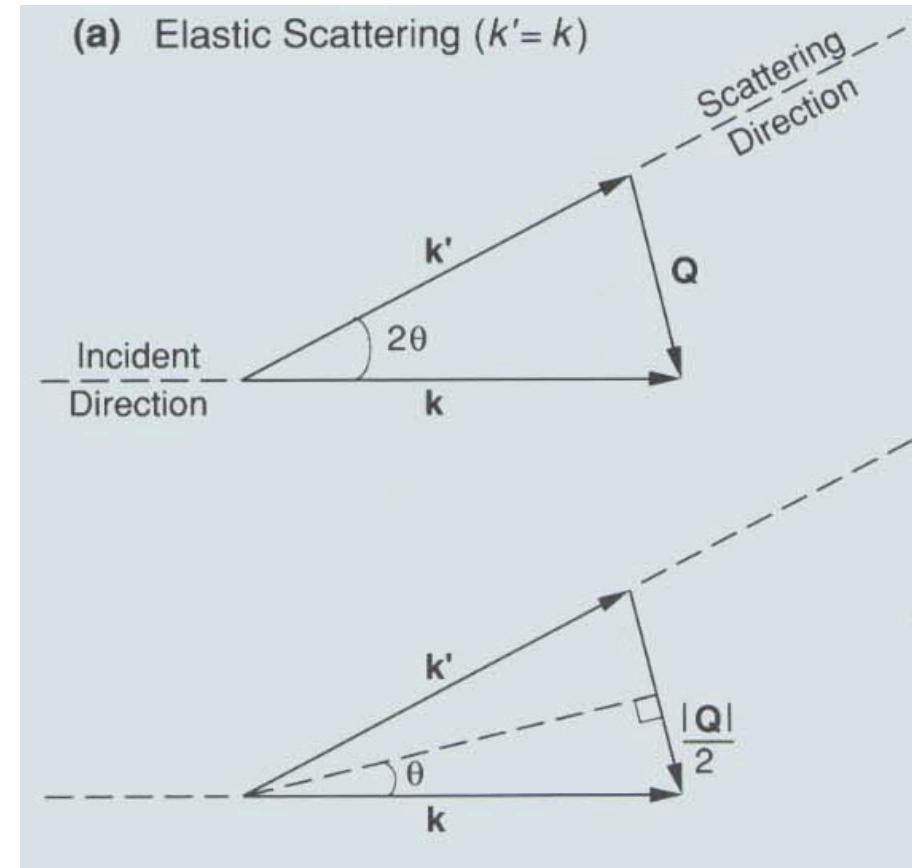
Scattering from single point



$$\sin \theta = \frac{q/2}{k}$$

$$q = 2k \sin \theta = \frac{4\pi \sin \theta}{\lambda}$$

Scattering vector



Scattering from a lattice

Scattering intensity

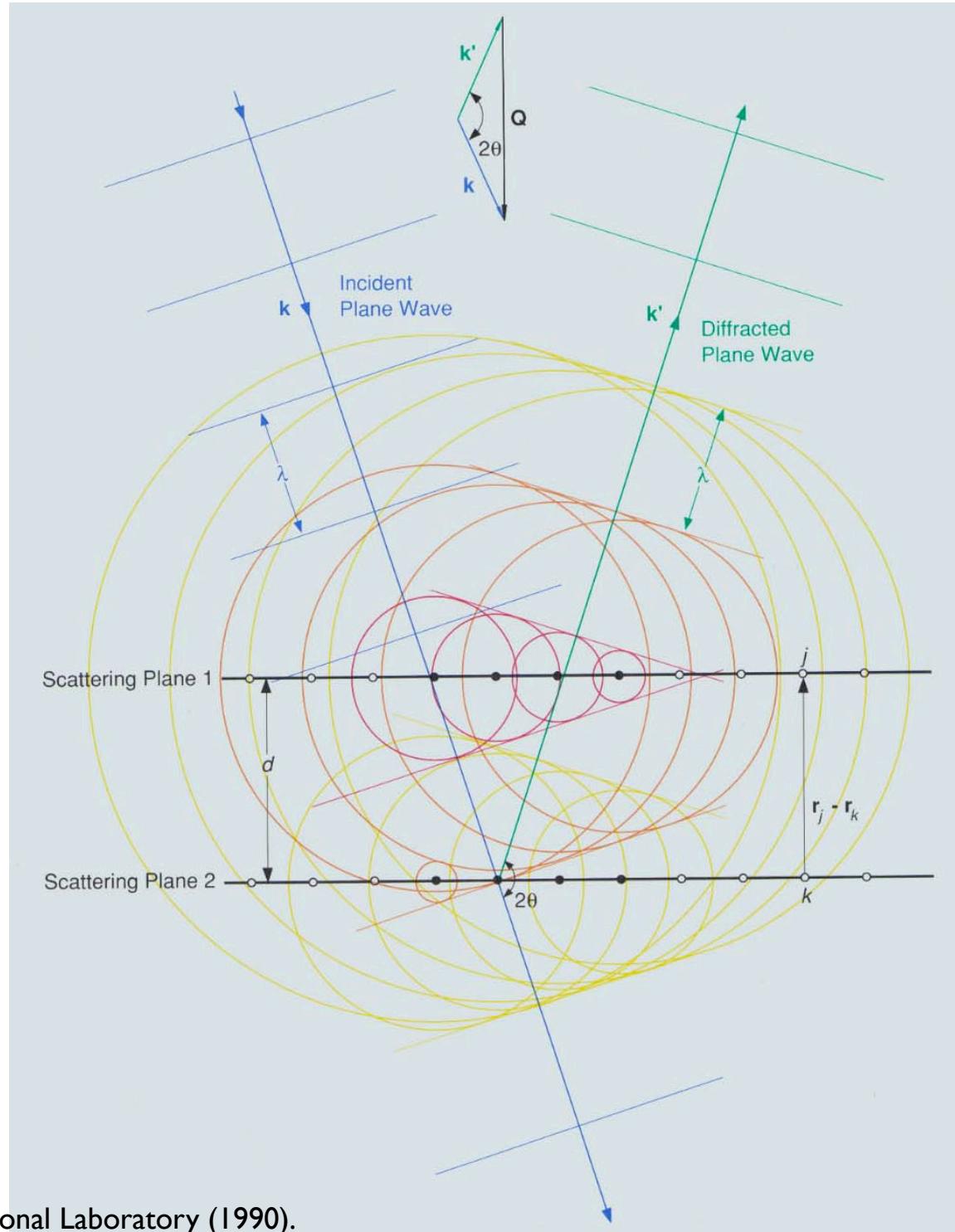
$$I(\mathbf{q}) = \sum_{j,k} b^2 e^{i\mathbf{q} \cdot (\mathbf{r}_j - \mathbf{r}_k)}$$
$$\equiv S(\mathbf{q})$$

Structure factor

$$\mathbf{q} \cdot (\mathbf{r}_j - \mathbf{r}_k) = qd = 2\pi n$$

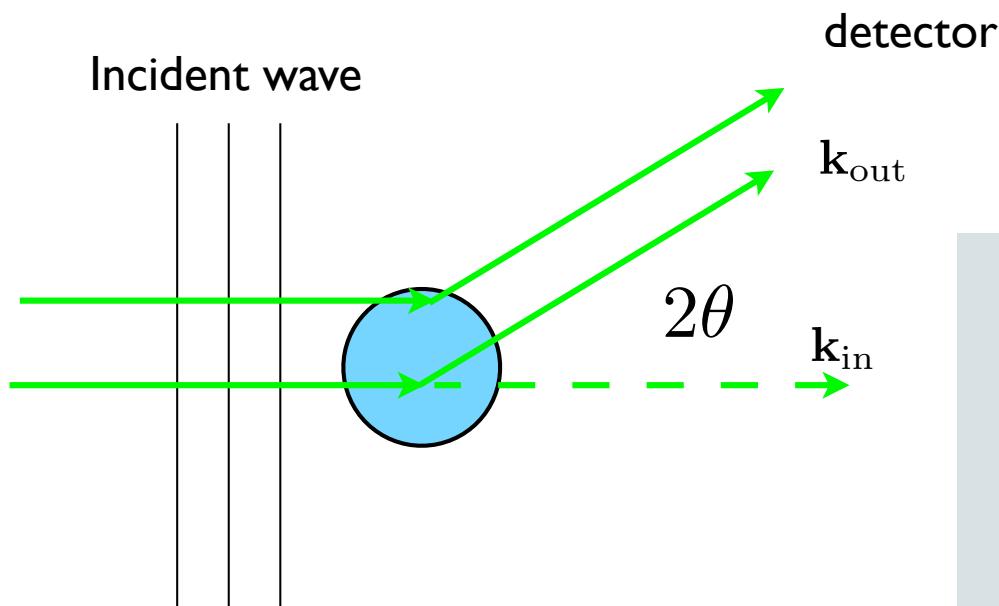
$$q = (4\pi/\lambda) \sin \theta$$

Bragg's law: $n\lambda = 2d \sin \theta$



Finite sized scatterer

Scattering from particle



$$P(q) = \left| \int_{V_p} d^3r e^{i\mathbf{q} \cdot \mathbf{r}} \right|^2$$

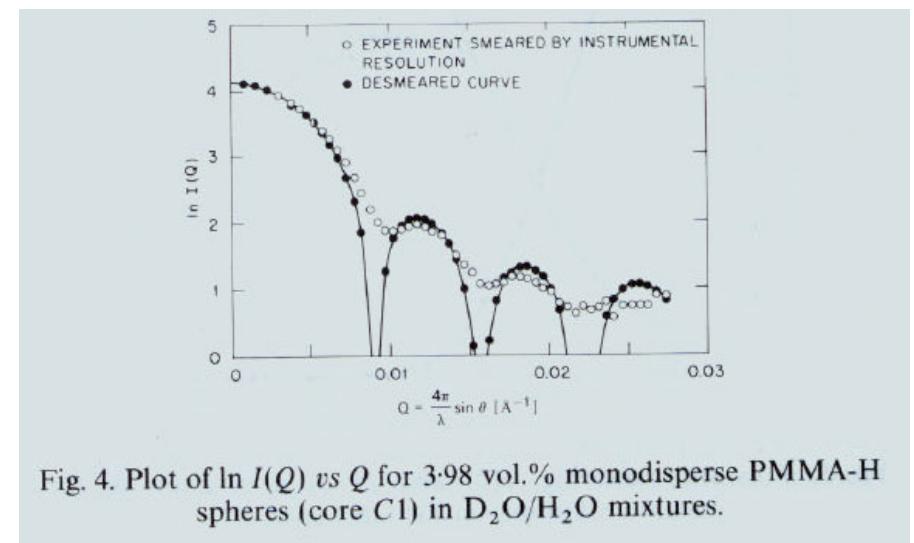
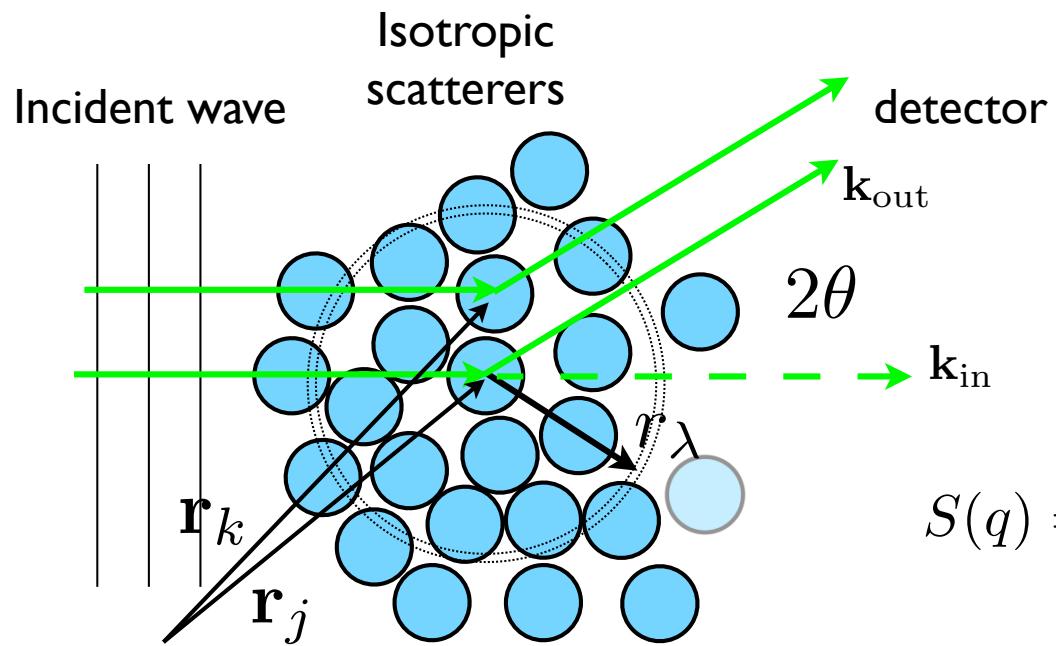


Fig. 4. Plot of $\ln I(Q)$ vs Q for 3.98 vol.% monodisperse PMMA-H spheres (core C1) in $\text{D}_2\text{O}/\text{H}_2\text{O}$ mixtures.

Total scattered intensity $I(q) \sim P(q)S(q)$

Scattering from random media



$$I(\mathbf{q}) = \sum_{j,k} b^2 e^{i\mathbf{q} \cdot (\mathbf{r}_j - \mathbf{r}_k)} \equiv S(\mathbf{q})$$

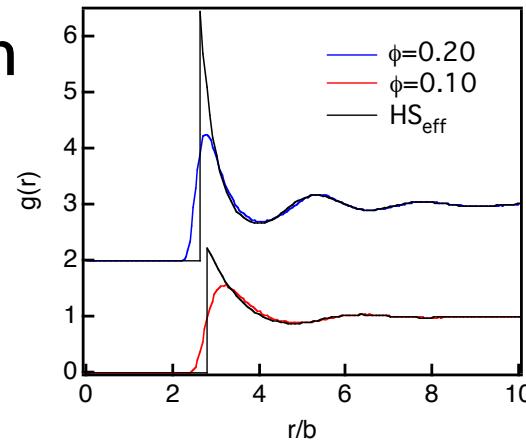
Structure factor relationship to RDF:

$$S(q) = 1 + \frac{4\pi n}{q} \int_0^\infty r(g - 1) \sin qr dr + \dots$$

Radial distribution function

$$g(r_\lambda) = \frac{3\langle N_\lambda \rangle}{2\pi N \rho(r_{\lambda+1}^3 - r_\lambda^3)}$$

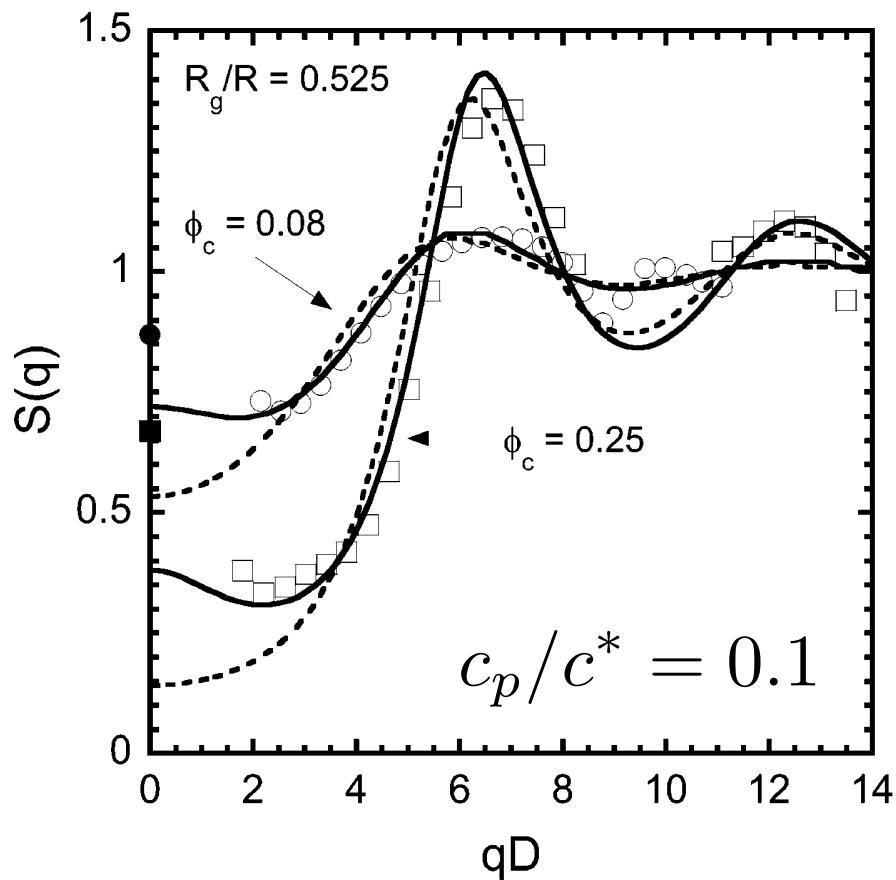
$$\langle N_\lambda \rangle \text{ in } r_\lambda < r < r_{\lambda+1}$$



Measure of average “local” structure

Structure factor

SANS, octadecyl coated silica



Relationship to RDF

$$S(q) = 1 + \frac{4\pi n}{q} \int_0^\infty r(g - 1) \sin qr dr + \dots$$

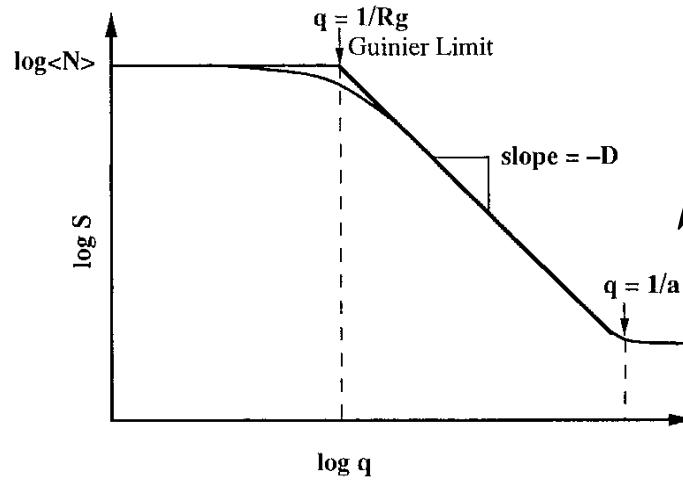
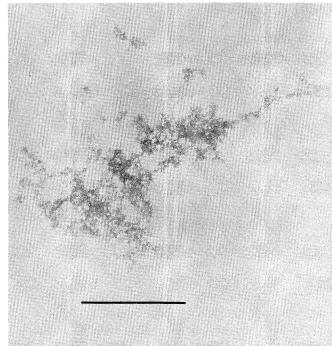
Isothermal compressibility

$$S(0, \phi) = \frac{d\phi}{d[\phi Z(\phi)]}$$

$$Z(\phi) = \Pi/nk_B T$$

Π Osmotic pressure

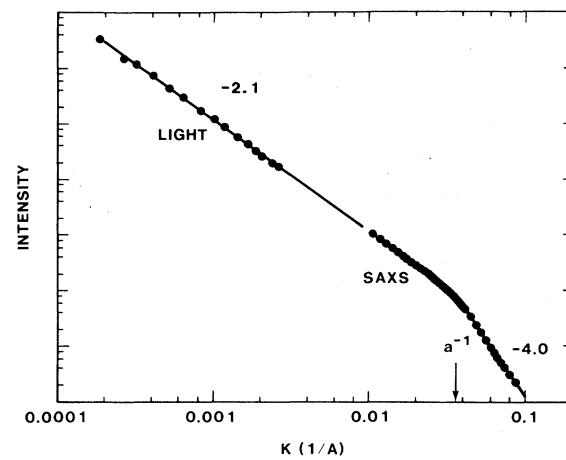
Fractal gels



$$S(q) \sim q^{-d_f}$$

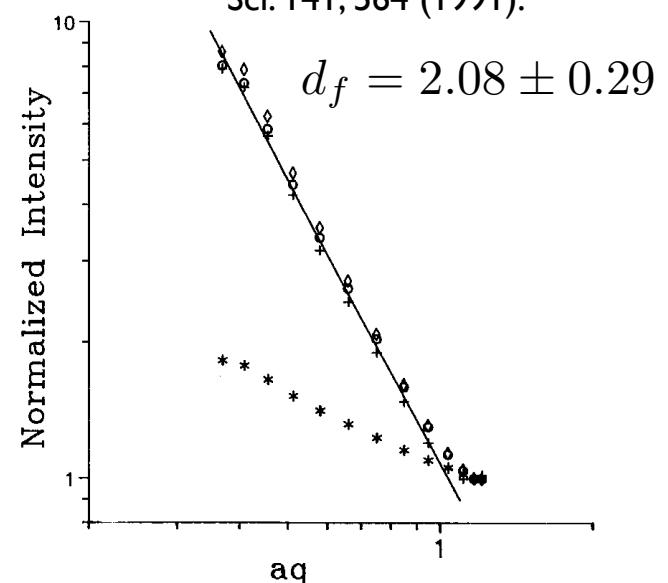
Ludox silica in 0.5M NaCl

D. W. Schaefer, et al., Phys. Rev. Lett. 52, 2371 (1984).



Stoeber silica in hexadecane

M. Chen and W. Russel, J. Colloid Interface Sci. 141, 564 (1991).

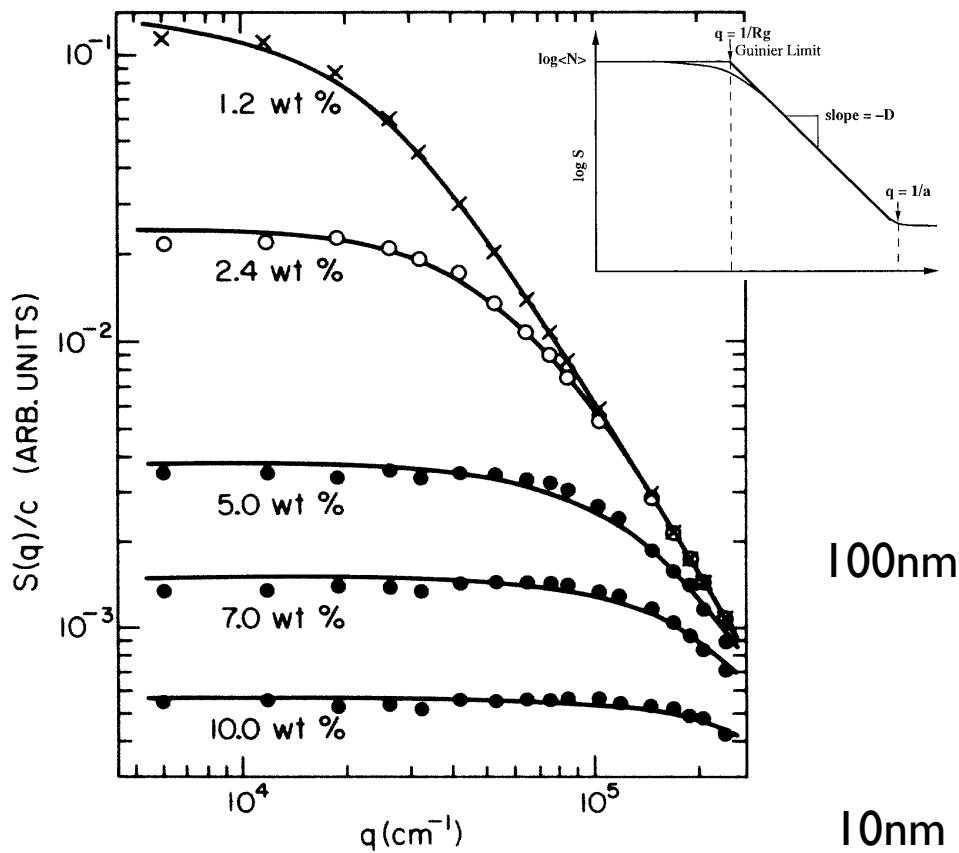


Fractal cut-off

G. Dietler, C. Aubert, D. S. Cannell, and P. Wiltzius, Phys. Rev. Lett. 57, 3117 (1986).

Structure factor

7nm colloidal silica
0.45M NaCl, pH 7.0

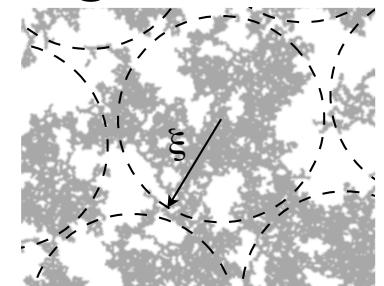


Model structure factor

$$S(q) = \frac{S(0)}{[1 + q^2 \xi^2 / (1 + \alpha)]^{1+\alpha}}$$

$$\alpha = 0.08 \quad d_f \approx 2.16$$

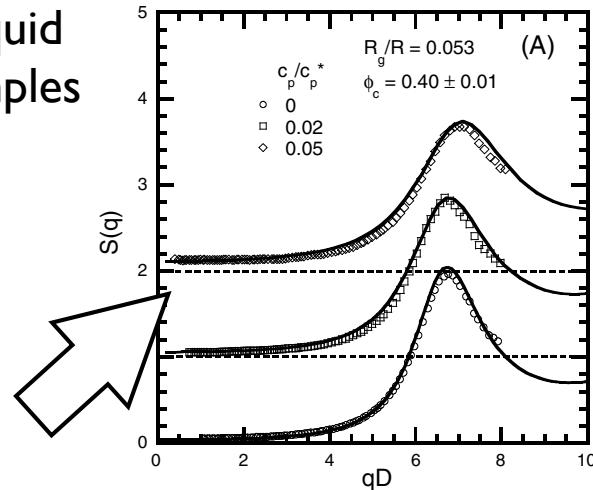
“Cross-over length”



Low-q structure

S.A. Shah, et al., J. Phys.: Condens. Matter 15, 4751 (2003).

Liquid samples



Gel samples

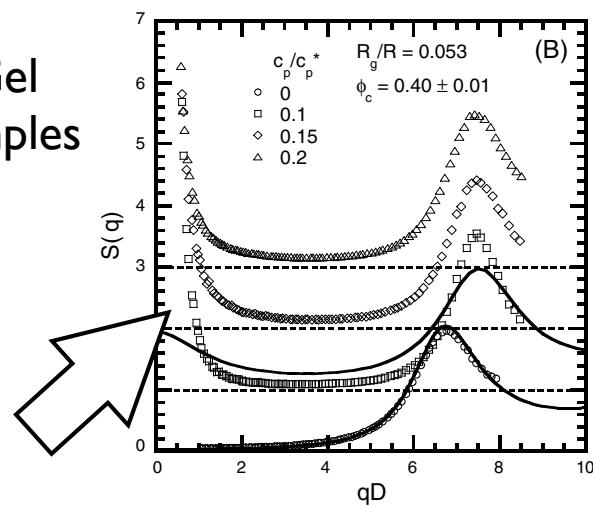


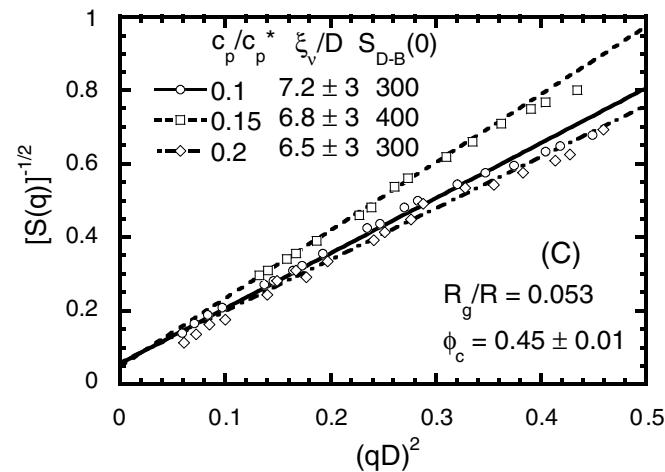
Figure 3. Structure factor $S(q)$ as a function of the dimensionless wavevector qD for samples with varying amounts of polymer quantified by the dimensionless polymer concentration, c_p/c_p^* . The colloid volume fraction is fixed at $\phi_c = 0.40 \pm 0.01$ and $R_g/R = 0.053$. (A) Liquid samples: $c_p/c_p^* = 0$ (circles), 0.02 (squares) and 0.05 (diamonds). (B) Gel samples: $c_p/c_p^* = 0$ (squares), 0.15 (diamonds) and 0.20 (triangles), and as a reference $c_p/c_p^* = 0$ (circles). Each subsequent data set is offset by 1 for clarity. Dotted lines show relevant baselines signifying $S(q) = 0$ for each data set. Full curves are the zero-adjustable parameter PRISM-mPY theory predictions.

- Low-q upturn: possibly void-like structures

Debye-Bueche random two-phase model:

$$S_{DB}(q) = \frac{S_{DB}(0)}{(1 + q^2 \xi_\nu^2)^2}$$

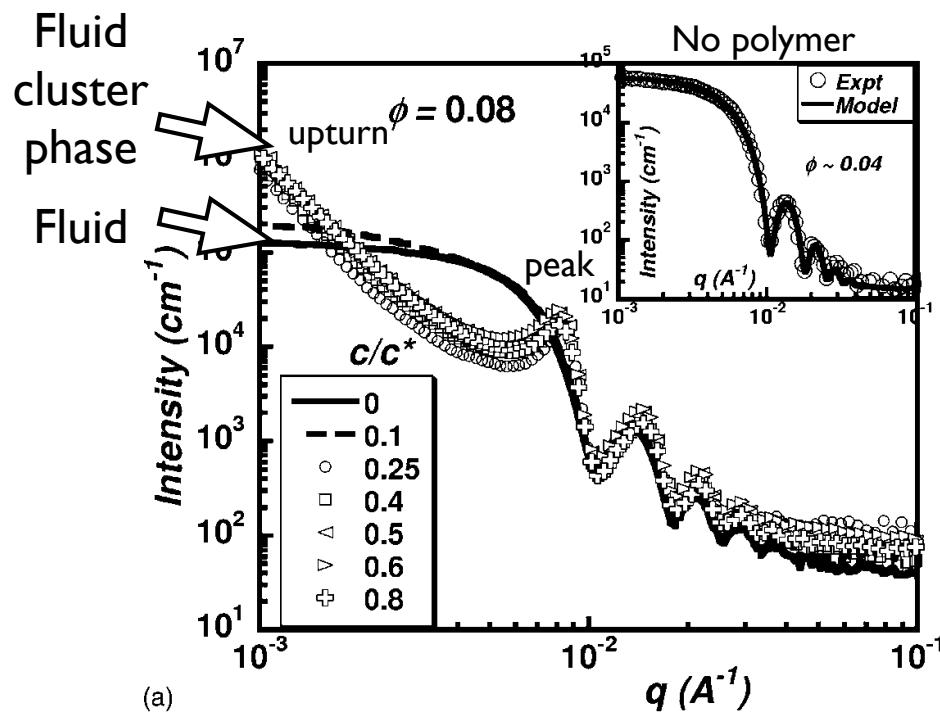
$$S_{DB}(0) = 48\phi \left(\frac{\xi_\nu}{D} \right)^3 \langle \eta^2 \rangle$$



Cluster precursors

V. Gopalakrishnan and C. F. Zukoski, Phys. Rev. E 75, 021406 (2007).

USAXS 42 nm silica, non-adsorbing polymer



- Mobile cluster phase precursor to gel.
- Microstructure of the particle clusters are similar to that observed in dense space-spanning depletion gels.
- Percolation of the clusters due to settling does not influence their internal microstructure.
- Supports hypothesis that formation of space-spanning depletion gels is driven by the percolation of particle clusters.

Microscopy

Basic imaging

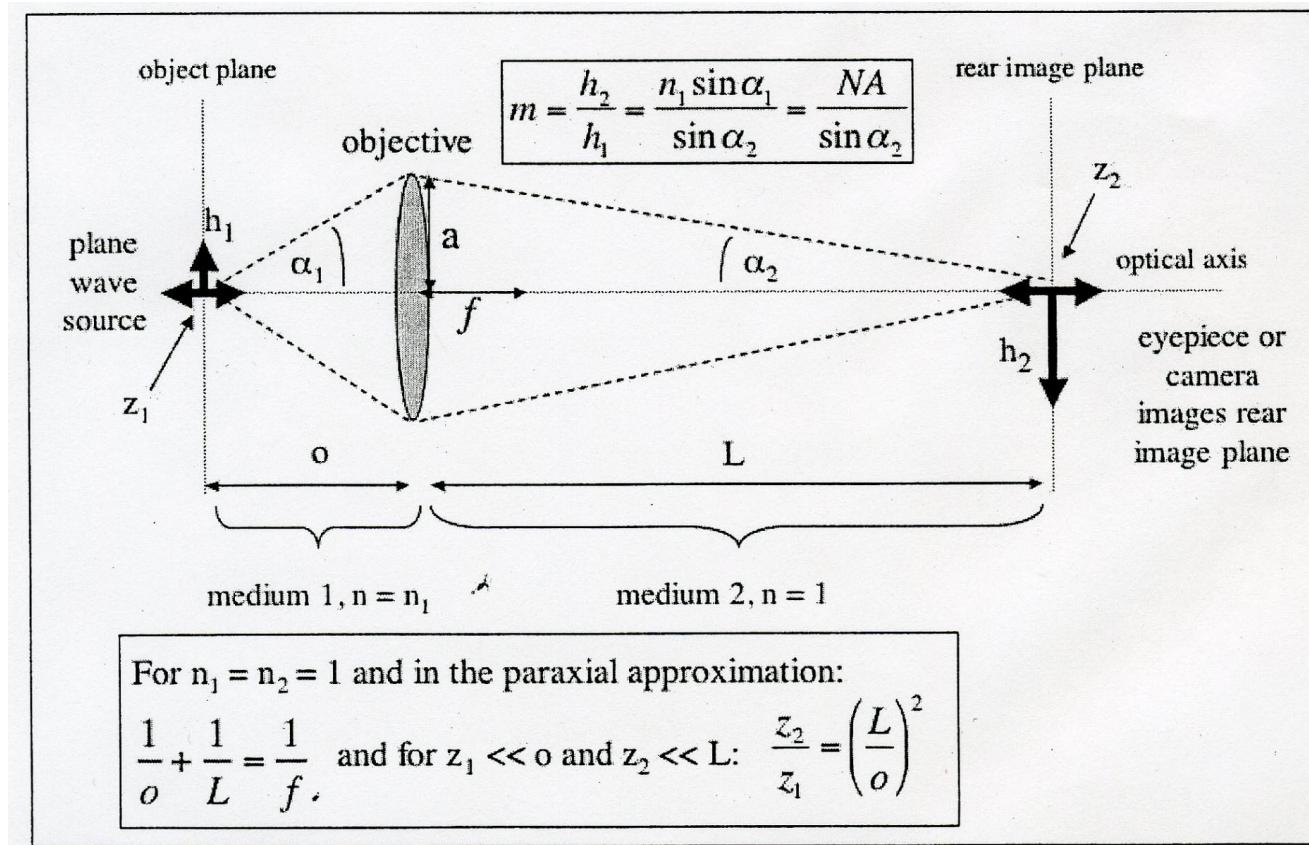
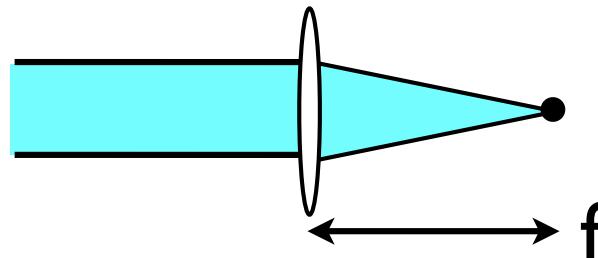
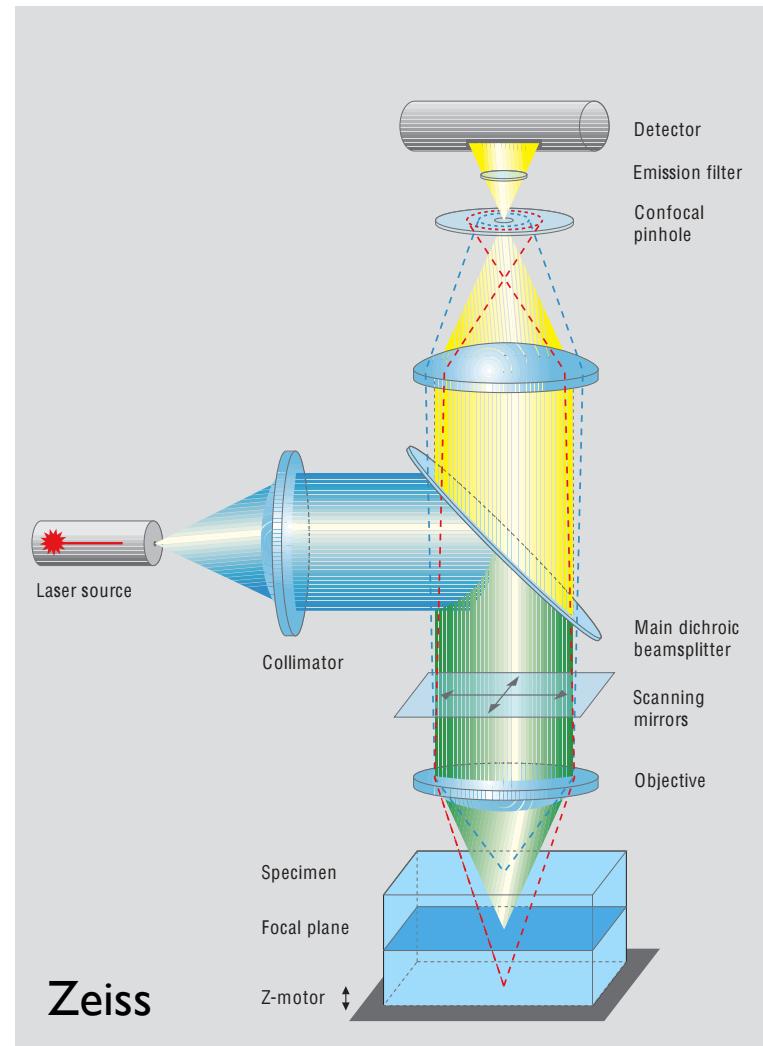


Image at infinity or
collimated beam



Confocal imaging

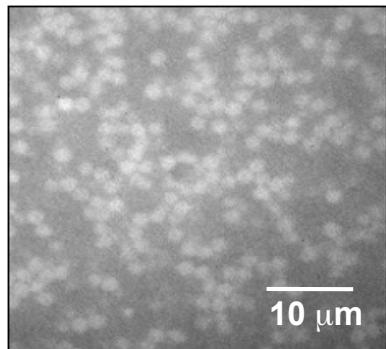
- Prevents out-of-plane rays from contributing to image.
- Useful for imaging dense suspensions.
- Objective or stage is moved to collect multiple, vertically stacked images.
- Stacks are combined to create image volumes.



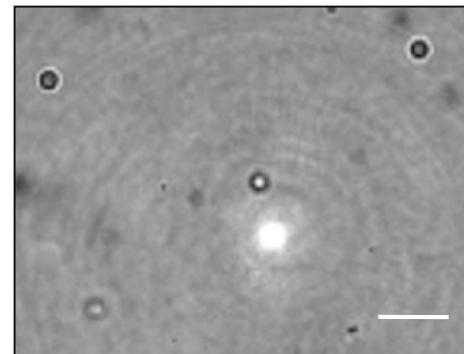
Confocal imaging in suspensions

- Requires specialized colloidal systems for index and density matching, i.e.

Fluorescent particle	Sterically-stabilized PMMA ($a = 0.77 \mu\text{m}$)
Index and density matched solvent	Decalin and cyclohexyl bromide
Non-adsorbing polymer	Polystyrene ($M_w = 900,000$), $c^* = 5\text{mg/ml}$



Fluorescent PMMA viewed with confocal microscope



Dispersed melamine particles viewed with bright-field microscope in same suspension

Particle tracking

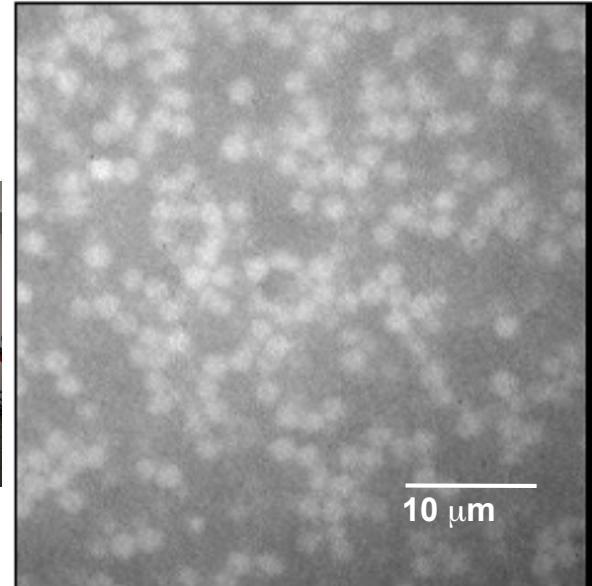
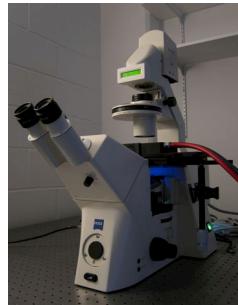
Crocker and Grier, J. Colloid Interface Sci., 179:298, 1996.

- Locate brightest pixel (x_0, y_0)

- Calculate weighted centroid

$$\begin{pmatrix} \varepsilon_x \\ \varepsilon_y \end{pmatrix} = \frac{1}{m_0} \sum_{i^2+j^2 \leq w^2} \begin{pmatrix} i \\ j \end{pmatrix} \cdot A(x+i, y+j)$$

$$m_0 = \sum_{i^2+j^2 \leq w^2} A(x+i, y+j)$$

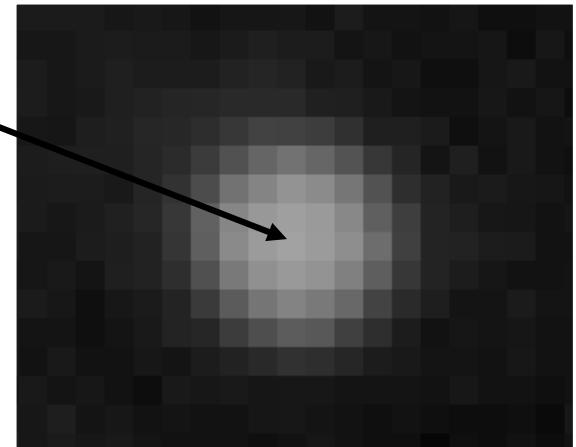


- Refine location

$$(x, y) = (x_0 + \varepsilon_x, y_0 + \varepsilon_y)$$

$$(x_0, y_0)$$

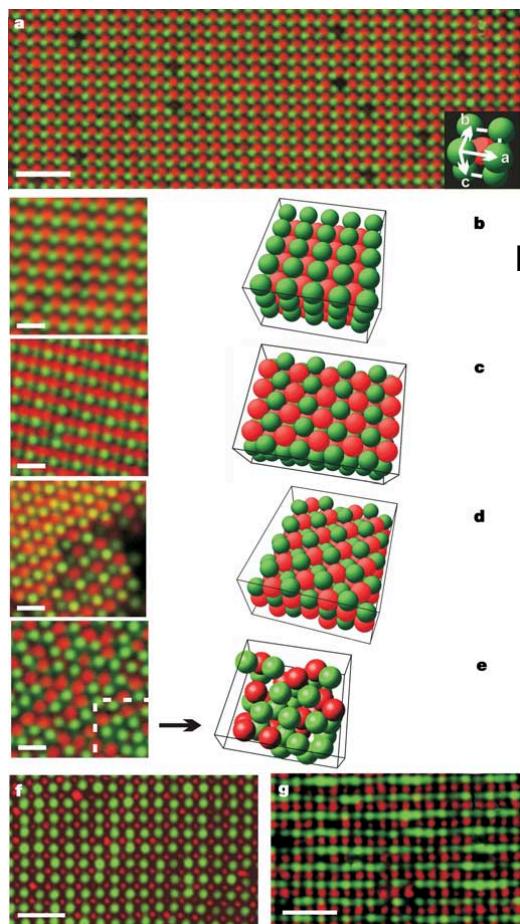
- Resolution to fraction of pixel (~10 nanometers)
- Correct static and dynamic errors
- See <http://www.physics.emory.edu/~weeks/idl/>
- Static and dynamic error:



T. Savin and P. S. Doyle, Biophys. J., 88:623–638, 2005.

Information: all particle positions in a 3D volume

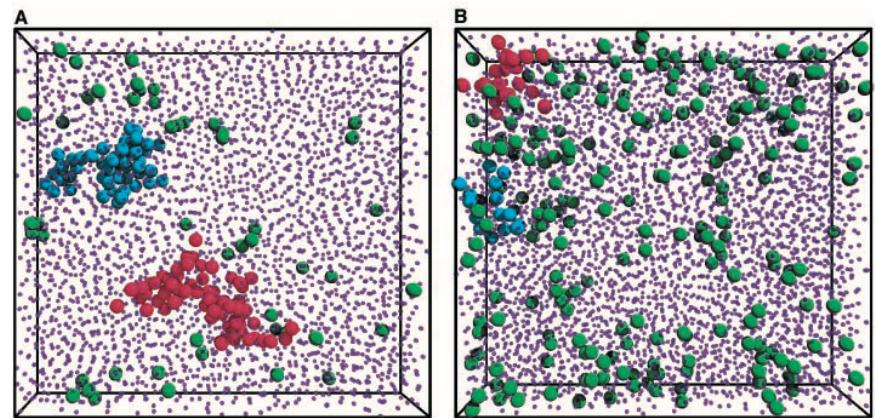
(Similar to simulations)



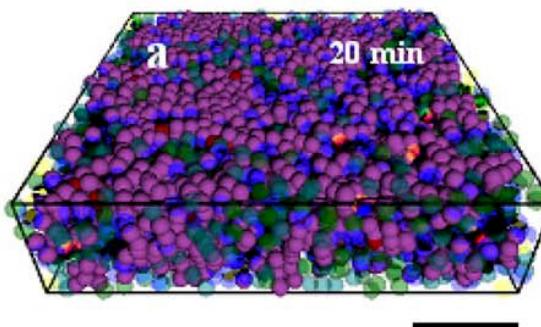
Colloidal crystals

M. E. Leunissen, et al., Nature 437, 236 (2005).

Glass dynamics



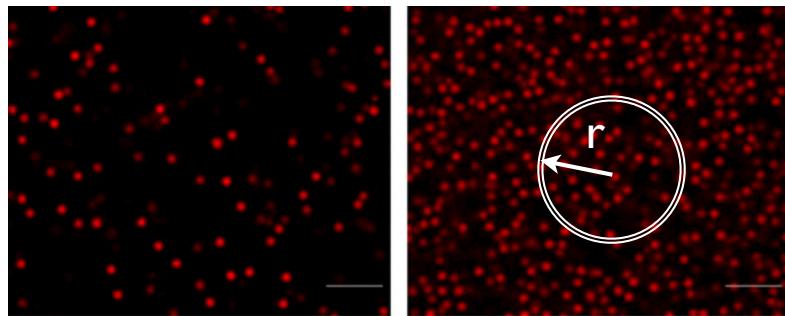
E. R. Weeks, et al., Science 287, 627 (2000).



Gels

Structure using microscopy

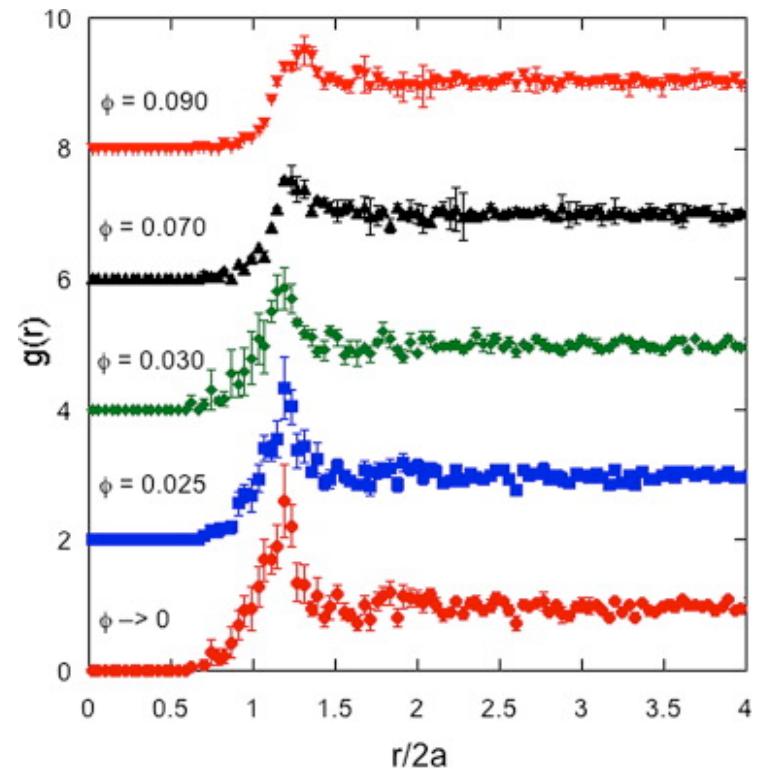
Radial distribution function



The radial distribution function is calculated directly from coordinates of particles.

$$g(r_\lambda) = \frac{3\langle N_\lambda \rangle}{2\pi N \rho (r_{\lambda+1}^3 - r_\lambda^3)}$$

$$\langle N_\lambda \rangle \text{ in } r_\lambda < r < r_{\lambda+1}$$



Images and data from M. Kogan, C. J. Dibble, R. E. Rogers, and M. J. Solomon, J. Coll. Int. Sci. 318, 252 (2007).

Integrated RDF

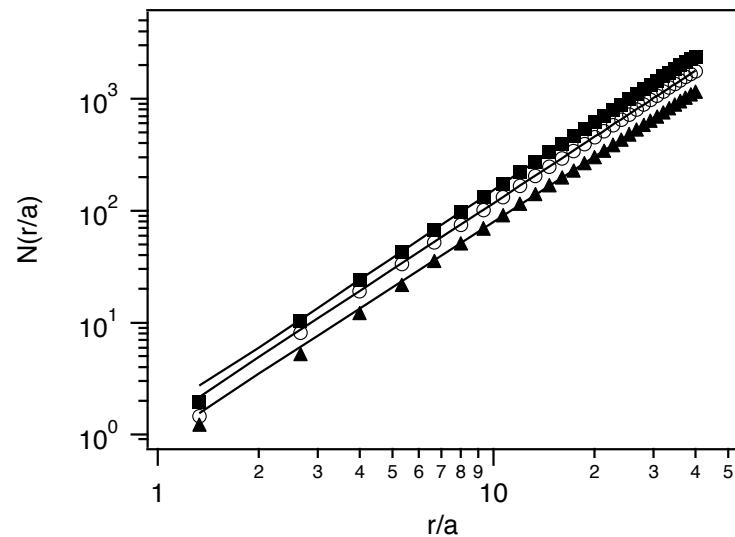
$$N(r) = 4\pi \int_0^r s^2 g(s) ds$$

Randomly distributed fractal clusters

$$g(r) \sim Ar^{d_f - 3} \exp(-r/\xi)$$

Power-law scaling of integrated RDF

$$N(r) = N_0 (\xi/a)^{d_f}$$



M. Hütter, J. Colloid Interface Sci. 231, 337 (2000).

F. Ferri, B. J. Friskin, and D. S. Cannell, Phys. Rev. Lett. 67, 3626, 1991.

Contact distribution

Local structure:
average number of
particles in “contact”

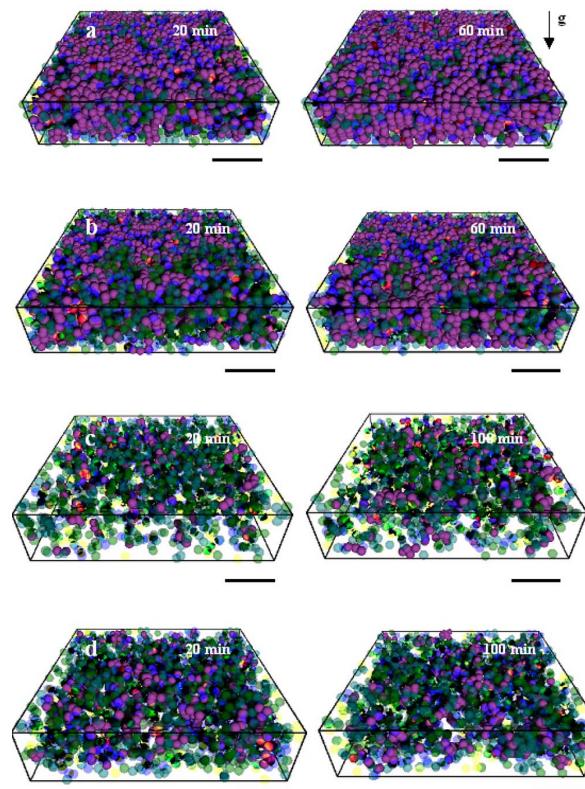
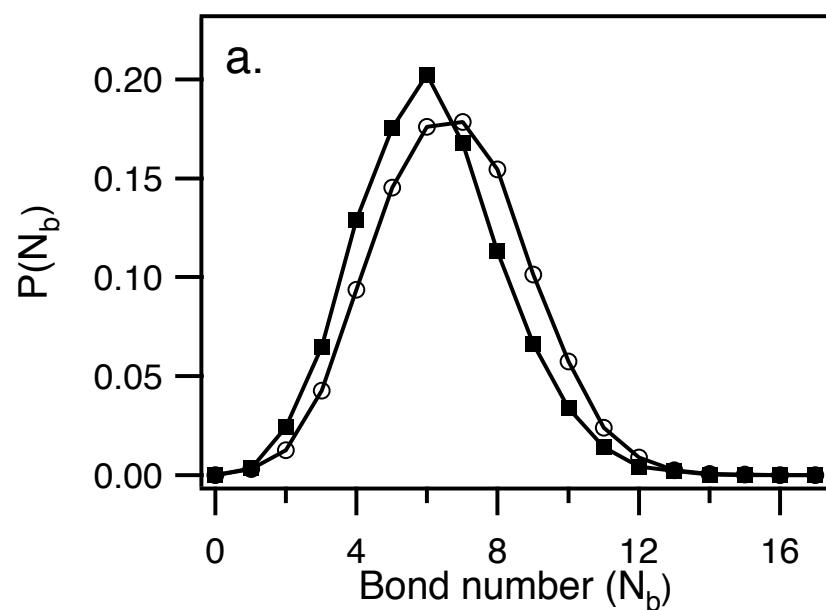


FIG. 9. Color online Visualizations of the structural evolution of sediments. a $C_p=0$ mg/ml; b $C_p=2.8$ mg/ml; c $C_p=4.1$ mg/ml; d $C_p=8.0$ mg/ml at $\phi=0.125$. Particles are described with seven different colors, depending the number of bonds per particles. A white circle means that a particle has no bond. Yellow, green, cyan, and blue represent one to four bonds, respectively. A pink indicates five to nine bonds. A red means greater than ten bonds. Particles with fewer than five bonds are rendered translucently.

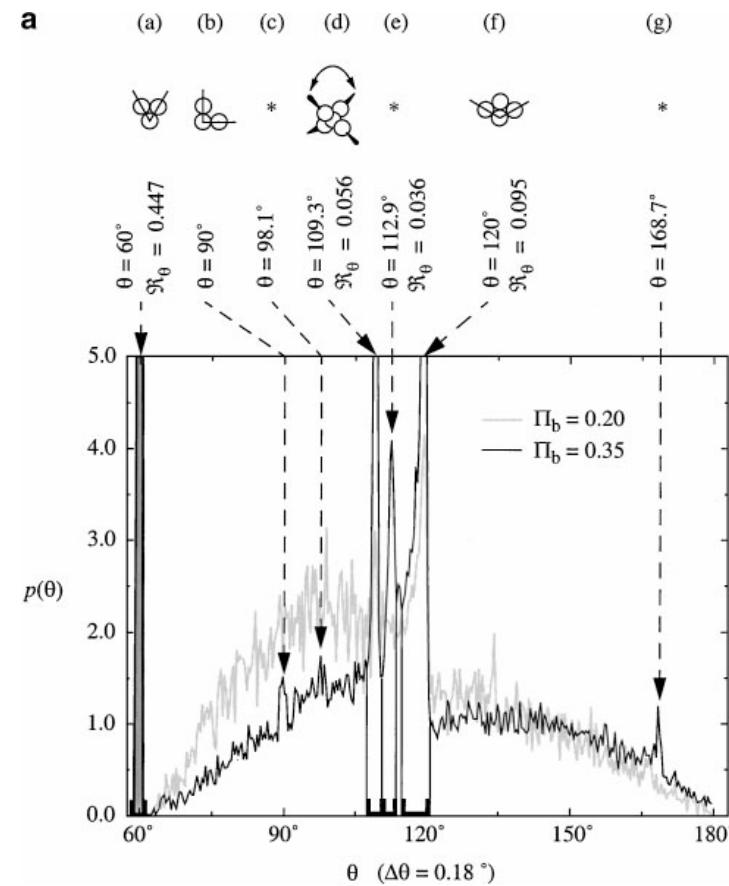
Bond angle distribution

$$p(\theta + \delta\theta/2) = \frac{2\pi}{3} \frac{\delta N_\theta(\theta, \theta + \delta\theta)}{\delta\theta N_\theta^t}$$

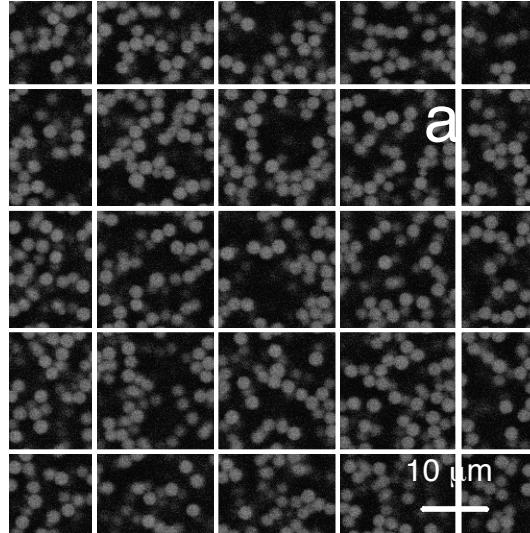
$\delta N_\theta(\theta, \theta + \delta\theta)$
particle triplets
with bond angle
between
 $\theta, \theta + \delta\theta$

N_θ^t

Total number
particle triplets



Density fluctuations

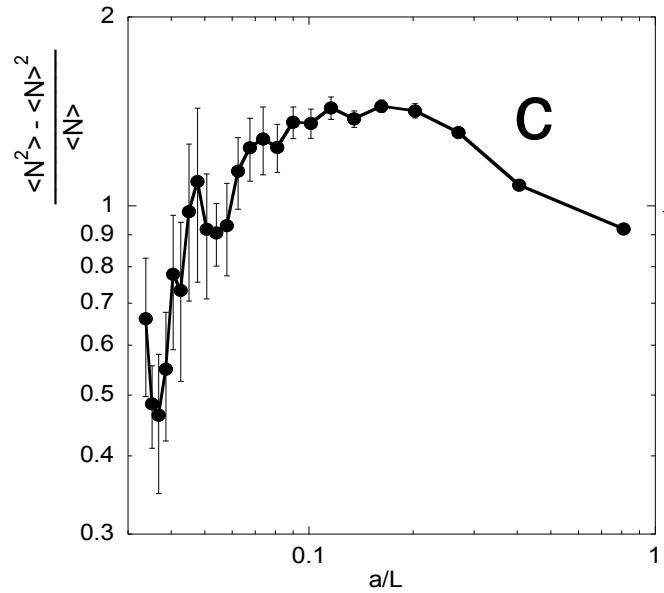


N Number of particles in volume L^3

$$\frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle^2}$$

Relationship to isothermal compressibility

$$\lim_{a/L \rightarrow 0} \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle^2} = S(0) = \rho k_B T \chi_T$$



M. Kogan, C. J. Dibble, R. E. Rogers, and M. J. Solomon, J. Coll. Int. Sci. 318, 252 (2007).

Higher concentrations

C. J. Dibble, M. Kogan, and M. J. Solomon, Phys. Rev. E 74, 041403 (2006)

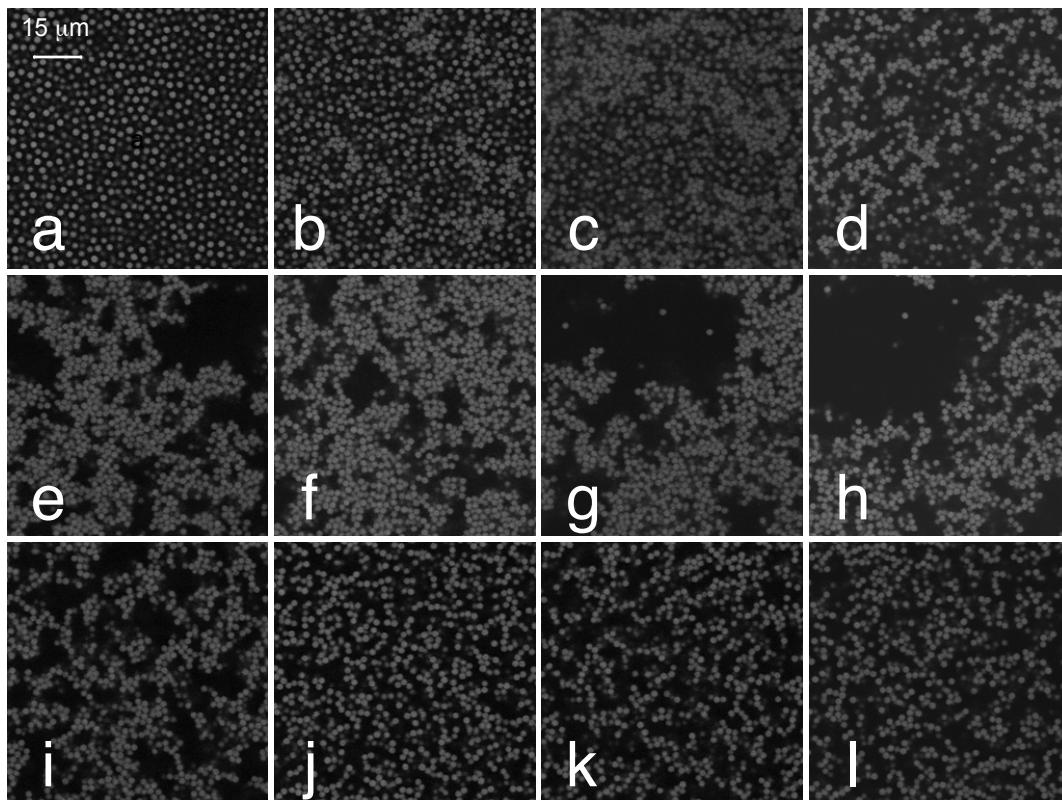
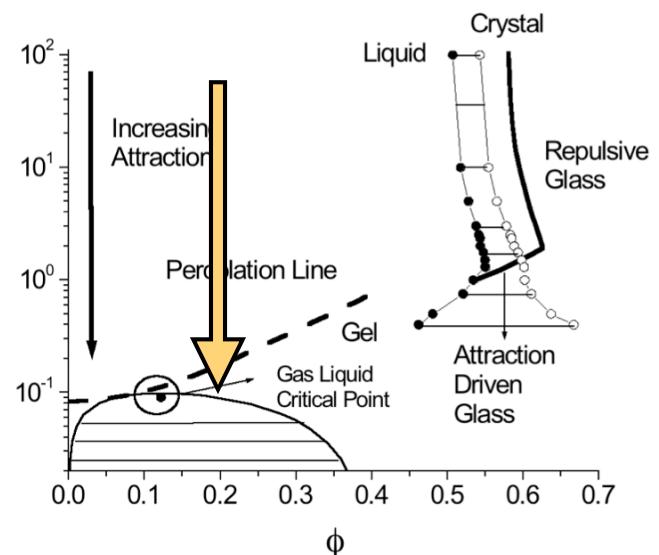


FIG. 3. Representative 2D confocal micrographs of system structure as a function of increasing depletion pair potential strength. (a) $c/c^* = 0.15$, (b) $c/c^* = 0.21$, (c) $c/c^* = 0.26$, (d) $c/c^* = 0.31$, (e) $c/c^* = 0.33$, (f) $c/c^* = 0.37$, (g) $c/c^* = 0.41$, (h) $c/c^* = 0.46$, (i) $c/c^* = 0.49$, (j) $c/c^* = 0.64$, (k) $c/c^* = 1.03$, (l) $c/c^* = 1.54$.



- a-free particles
- b-mobile clusters coexist with free particles
- c-increasingly immobile clusters
- d,e,f,g-bicontinuous network
- h-maximum void structure
- i-more homogeneous, branched
- j,k,l-chain-like locally, homogeneous globally

Density fluctuations

C. J. Dibble, M. Kogan, and M. J. Solomon, Phys. Rev. E 74, 041403 (2006).

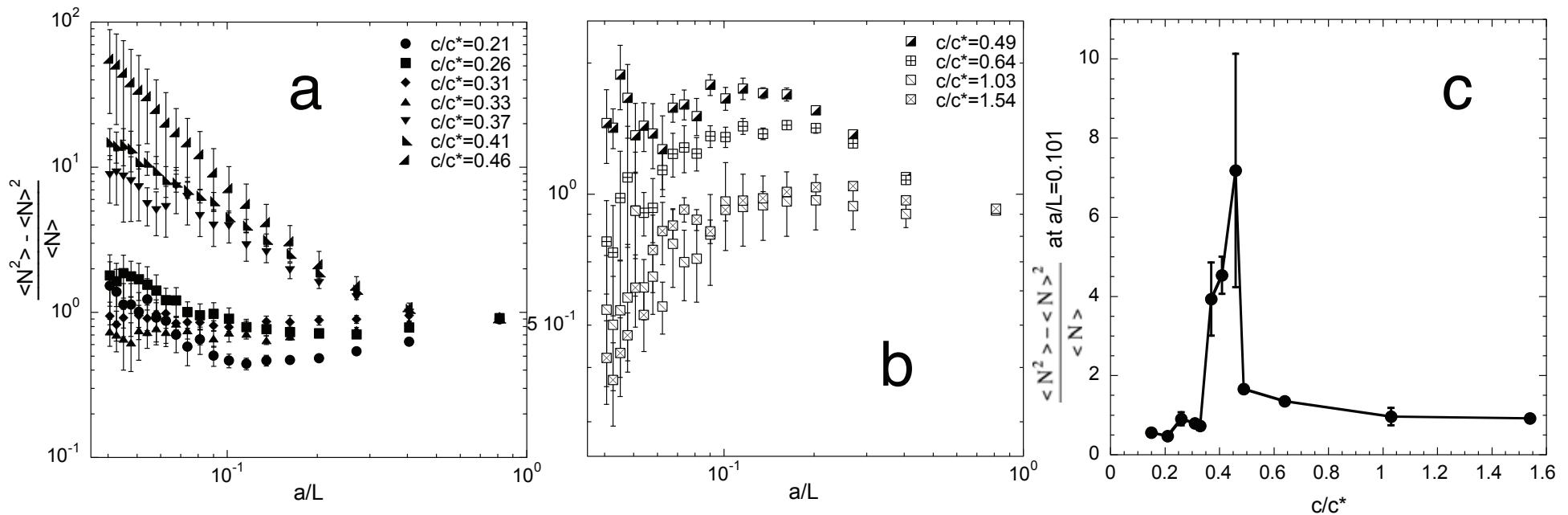


FIG. 6. Number density fluctuation measures show the normalized mean-squared fluctuations of particle number as a function of subsystem size considered as it varies with attraction strength for a $c/c^* = 0.46$ and b $c/c^* = 0.49$. c Mean number-density fluctuations for all attraction strengths for inverse length scale $a/L=0.101$.

Contact number distribution

C. J. Dibble, M. Kogan, and M. J. Solomon, Phys. Rev. E 74, 041403 (2006).

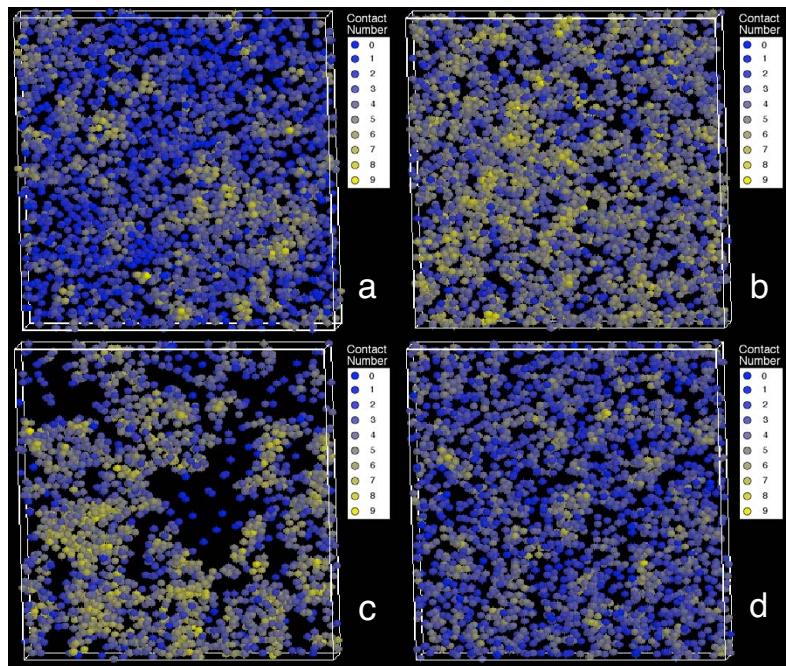
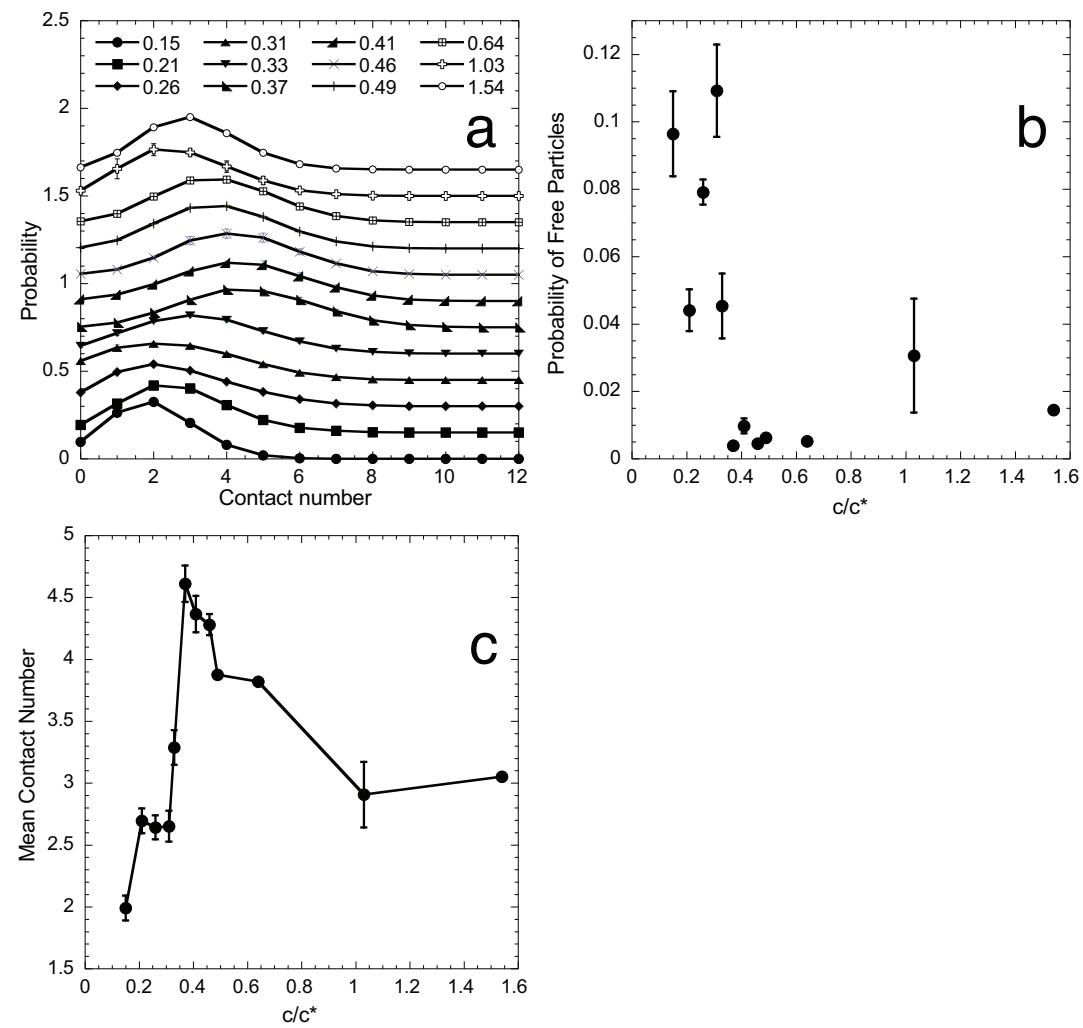


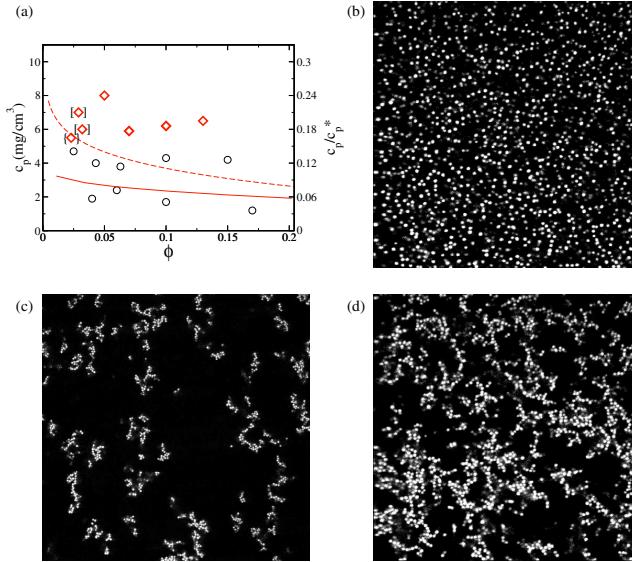
FIG. 4. (Color online) Shaded representation depicts spatial dependence of contact number on depletant concentration for (a) $c/c^* = 0.26$, (b) $c/c^* = 0.37$, (c) $c/c^* = 0.41$, (d) $c/c^* = 1.03$ where lighter shading (yellow color online) correspond to high contact numbers and darker shading (blue color online) correspond to low contact numbers.



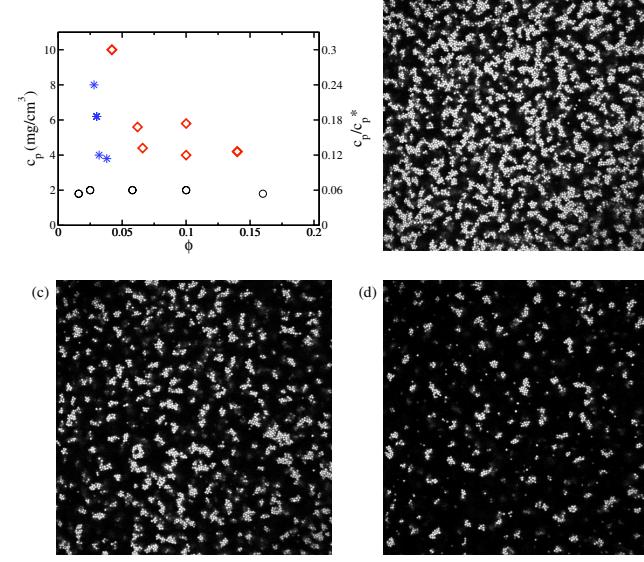
“Route to gelation is highly dependent on the state of density matching and charge of the constituent particles”

H. Sedwick, S. U. Egelhaaf, and W. C. K. Poon, J. Phys: Condens. Matter 16, S4913 (2004).

Screened, density-matched



Screened, density-mismatched



Charged, density-matched

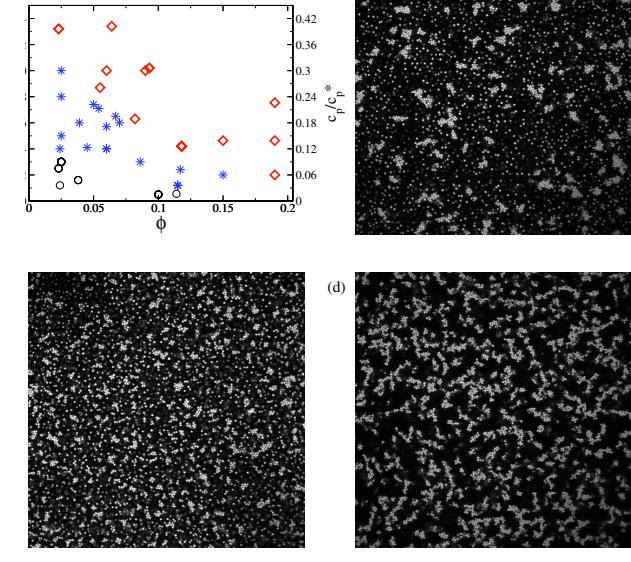


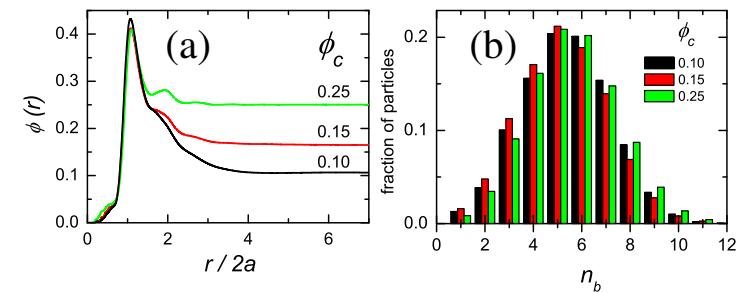
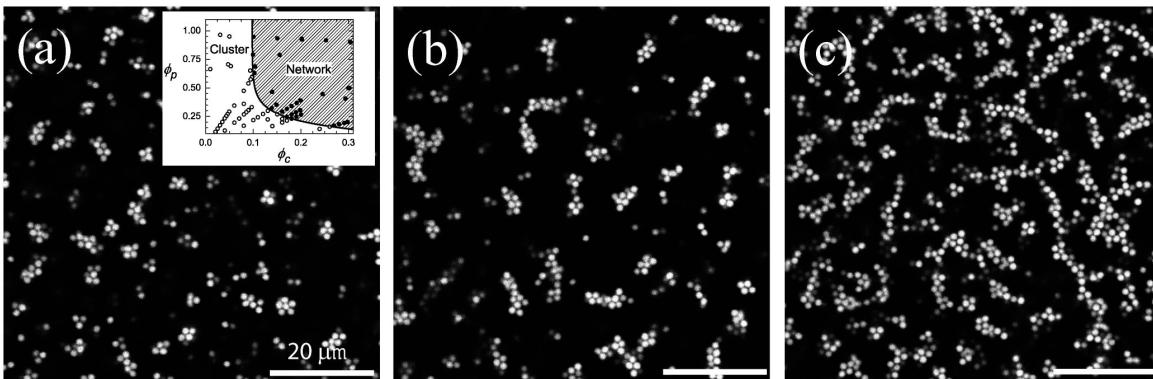
Figure 3. (a) Observations of the charged, density-matched system ($Q\bar{G}$) of $\delta \approx 0.021$. Symbols represent the following: fluid—circle, mobile clusters—star, and frozen structures—diamond. (b) A confocal micrograph of the clusters in a sample with $\phi = 0.086$ and $c_p = 3 \text{ mg cm}^{-3}$. (c) A confocal micrograph of the clusters in a sample with higher volume fraction, $\phi = 0.15$, and $c_p = 1.9 \text{ mg cm}^{-3}$. (d) A confocal micrograph of the frozen structure at $\phi = 0.083$ and $c_p = 6 \text{ mg cm}^{-3}$.

Long-range repulsion

A. I. Campbell, V. J. Anderson, J. S. van Duijvenveldt, and P. Bartlett,
Phys. Rev. Lett. 94, 208301 (2005).

Transition from cluster phase to gel

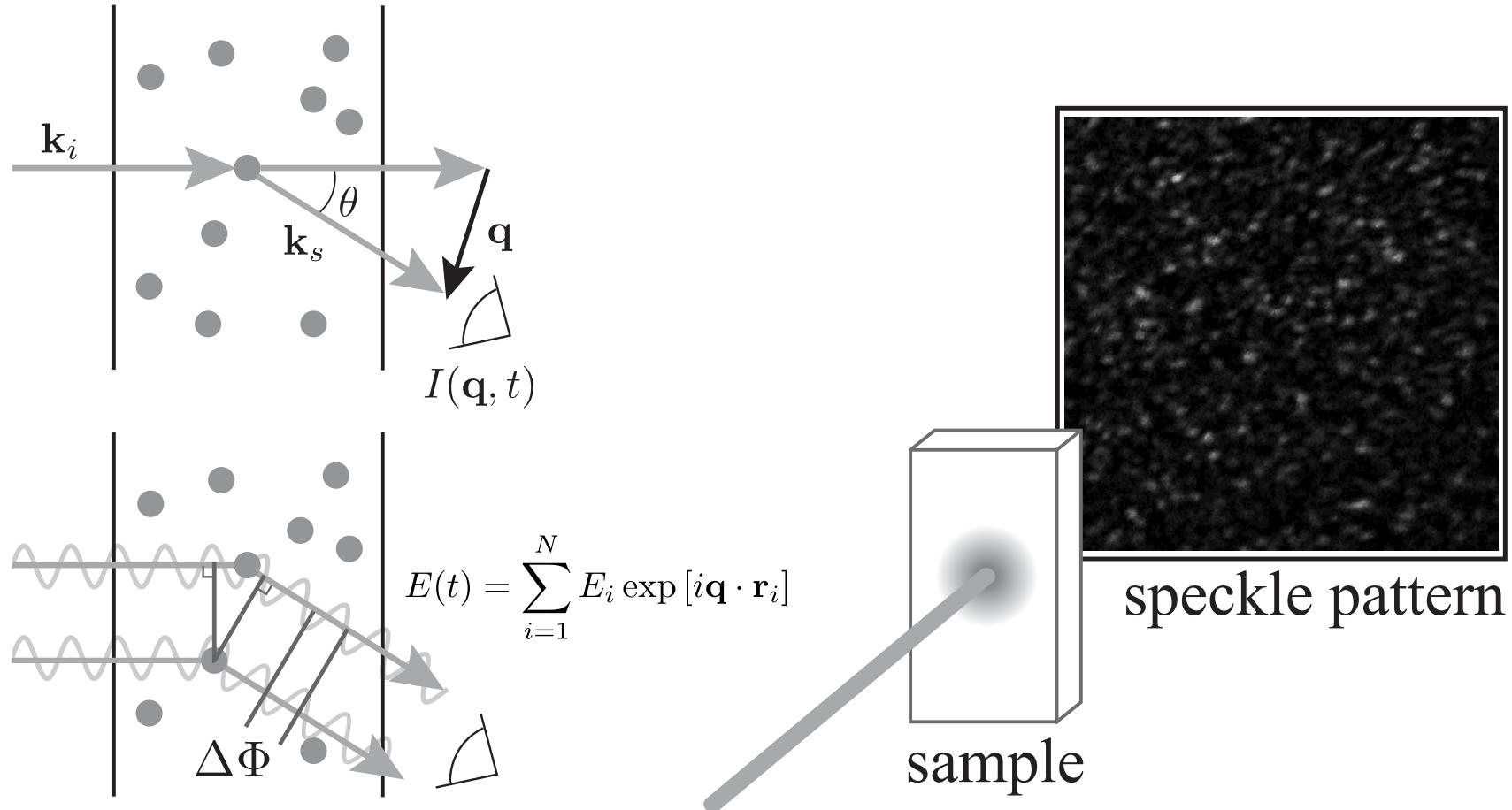
Unusual local order in network:
face-sharing tetrahedral clusters



- Long-range repulsive interactions alters dynamical arrest.
- Formation of stable, chain-like clusters, constructed from face-sharing tetrahedral units, which associate into a fully connected network at gelation.
- Does formation of tetrahedra reflect the presence of charge or is simply a consequence of a centrosymmetric attractive potential?

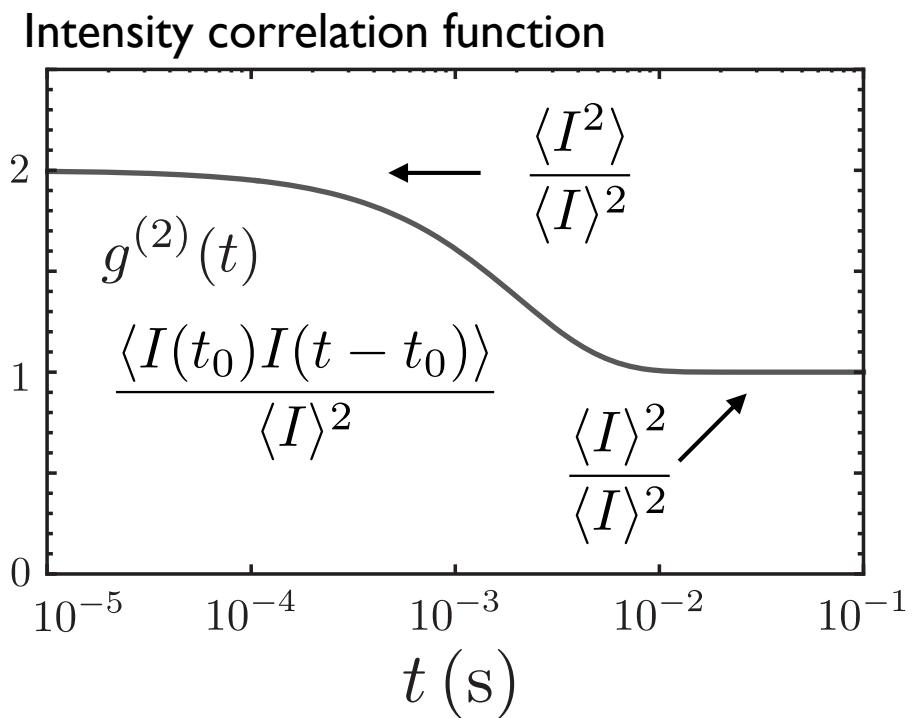
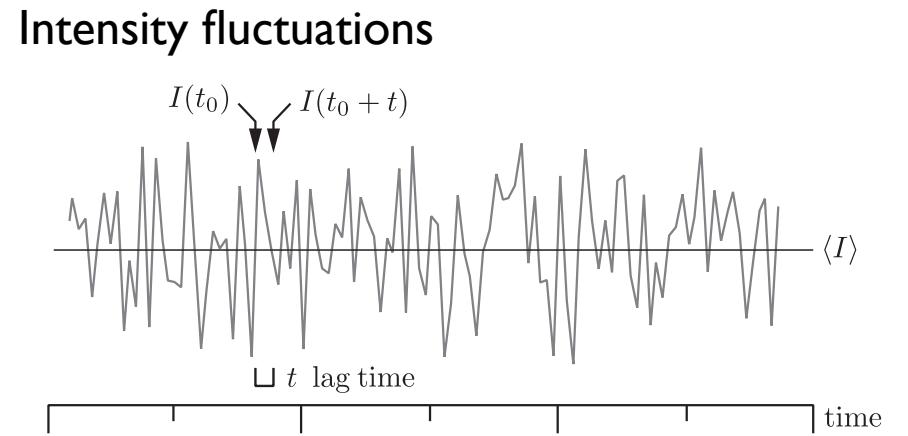
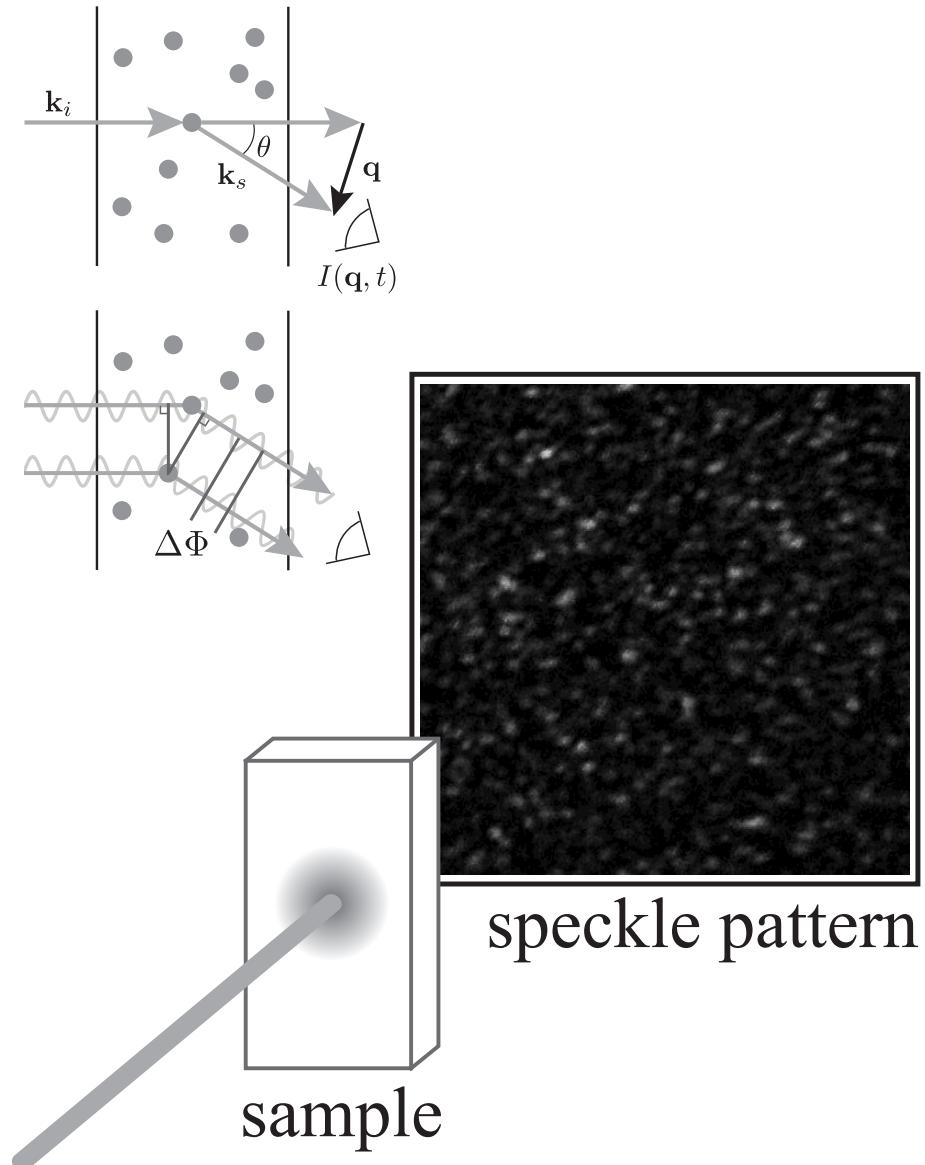
Dynamic scattering

Light scattering



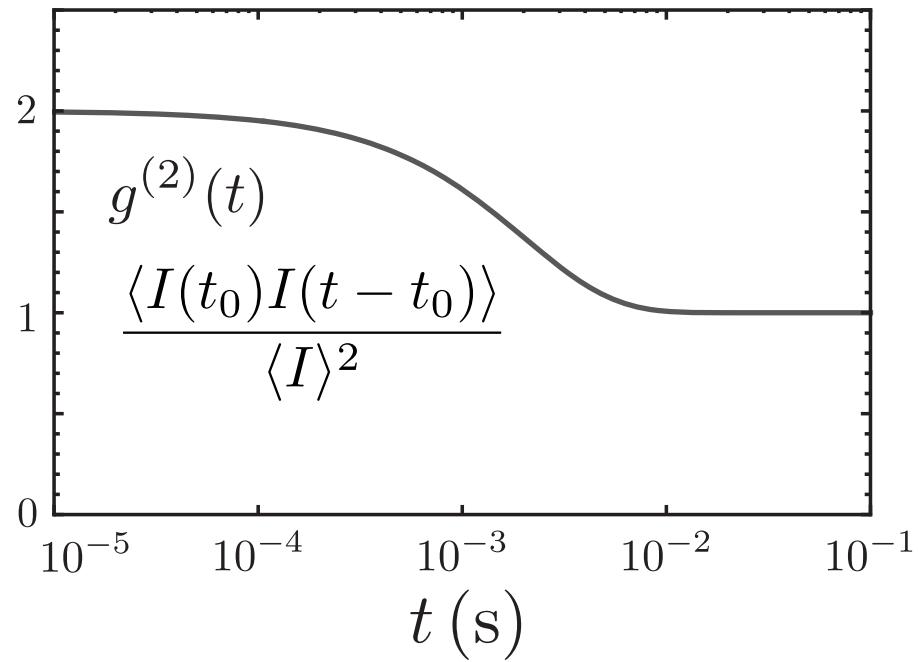
$$q = |\mathbf{q}| = (4\pi n / \lambda_v) \sin \theta / 2 \quad \text{scattering vector}$$

Intensity fluctuations



Light scattering microrheology

Intensity correlation function



Field correlation function

$$g^{(1)}(t) = \frac{\langle E(t)E^*(0) \rangle}{\langle |E(t)|^2 \rangle}$$

$$g^{(1)}(t) = \exp \left[\frac{-q^2 \langle \Delta r^2(t) \rangle}{6} \right]$$

Siegert relation

$$g^{(2)}(t) = 1 + \beta |g^{(1)}(t)|^2$$

Internal dynamics and elasticity of fractal colloidal gels

Krall, A. H. & Weitz, D. A. Phys. Rev. Lett. 80, 778–781 (1998).

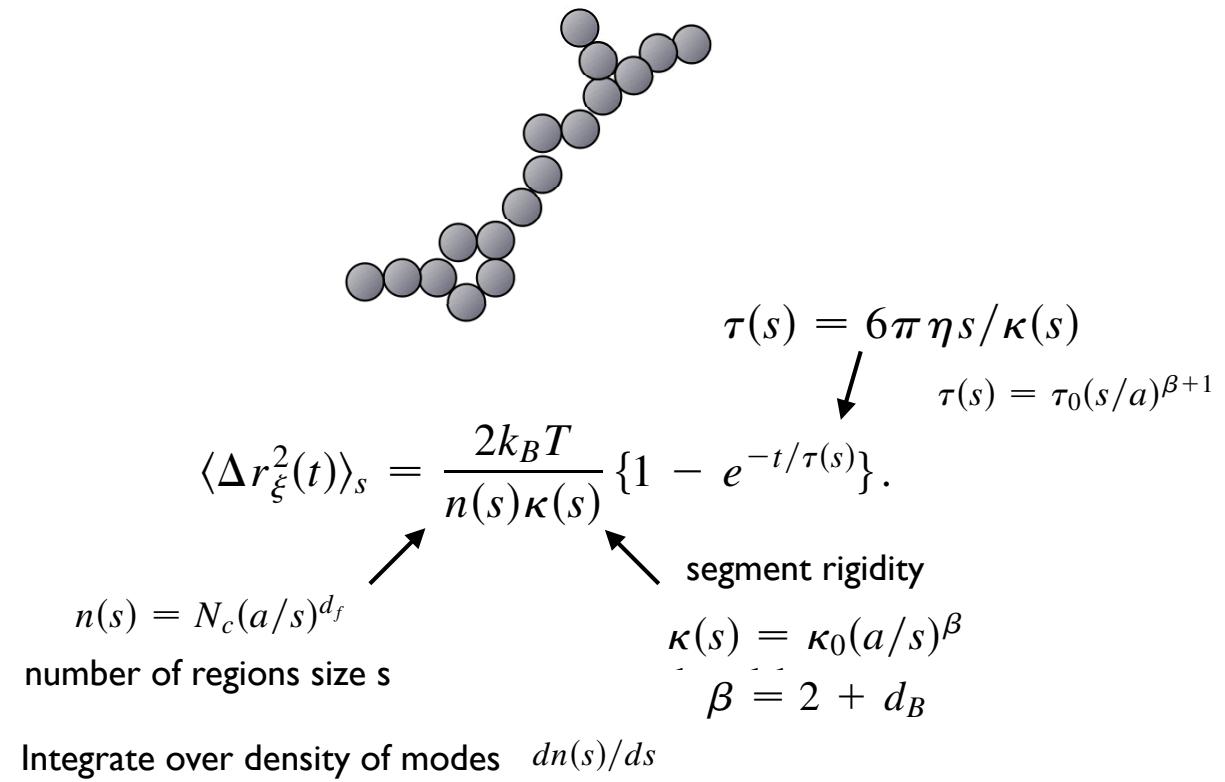
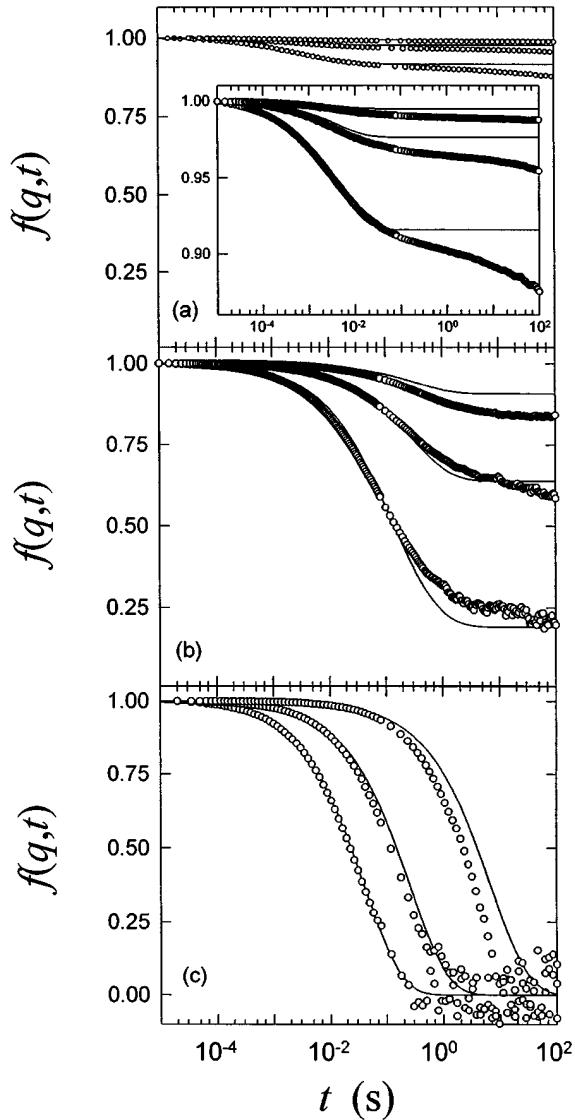


FIG. 1. Dynamic structure factors $f(q,t)$ of fractal colloidal gels for different ϕ_0 and q ; (a) $\phi_0 = 5.0 \times 10^{-3}$, and $q = 4.1, 8.7$, and $16.7 \mu\text{m}^{-1}$ from top to bottom; smooth curves are theoretical fits yielding $\delta^2 = 1.9 \times 10^{-3} \mu\text{m}^2$ and $\tau = 3.9 \times 10^{-3}$ s; (b) $\phi_0 = 1.5 \times 10^{-3}$; $q = 4.1, 8.7$, and $16.7 \mu\text{m}^{-1}$; $\delta^2 = 3.5 \times 10^{-2} \mu\text{m}^2$, $\tau = 0.36$ s; and (c) $\phi_0 = 1.7 \times 10^{-4}$; $q = 4.1, 12.1$, and $22.3 \mu\text{m}^{-1}$; $\delta^2 = 13 \mu\text{m}^2$, $\tau = 1850$ s.

Methods summary

Technique	Advantages	Disadvantages
Neutron	<ul style="list-style-type: none">• Works for many materials• Bulk samples (1-2mm)• Rheo-SANS	<ul style="list-style-type: none">• Ensemble averaged structure• Accessibility
X-ray	<ul style="list-style-type: none">• Works for many materials	<ul style="list-style-type: none">• Ensemble averaged structure• Shorter penetration lengths• Accessibility
Light	<ul style="list-style-type: none">• Convenience• Simultaneous rheology	<ul style="list-style-type: none">• Ensemble averaged• Specialized experimental systems (index match)
Confocal	<ul style="list-style-type: none">• Direct, real space structure• Simultaneous dynamics	<ul style="list-style-type: none">• Specialized experimental systems (index, density match)• Flow studies limited by frame rate
Electron	<ul style="list-style-type: none">• Detailed particle size, shape information	<ul style="list-style-type: none">• Sample preparation• Vacuum technique